

ITMP seminar, Zoom, May 2025

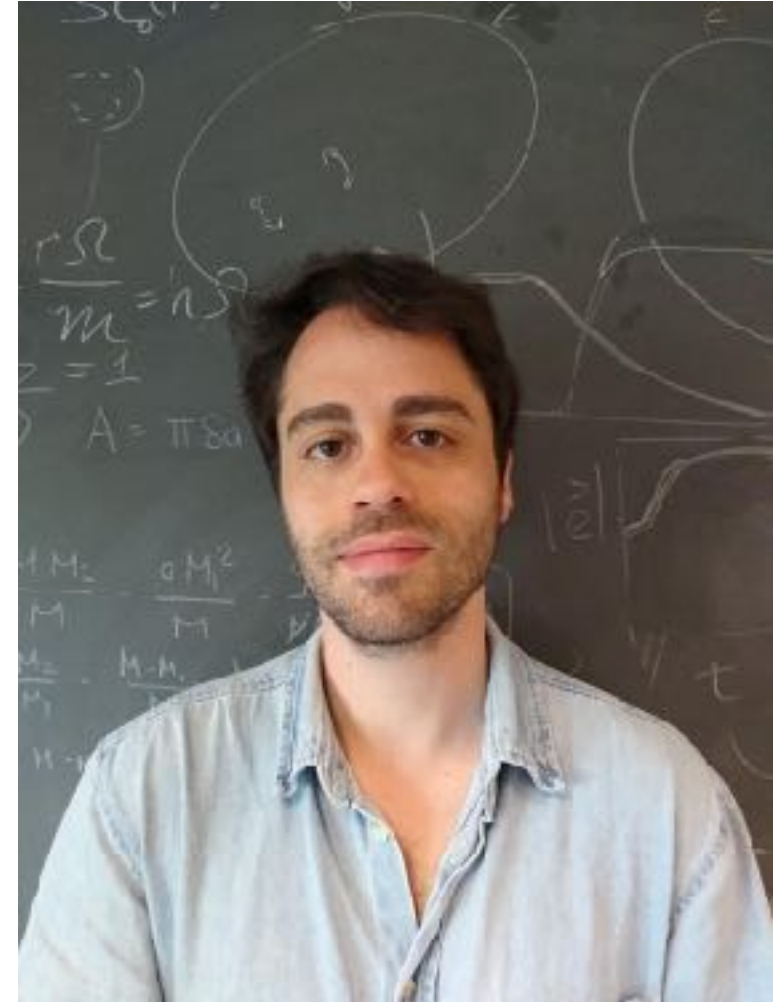
Amplitudes for Hawking Radiation

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With



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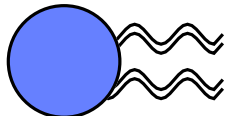
Motivation

Motivation

Classical gravitational waveform: from (generalised) amplitudes:

$$\text{waveform}(t) = \int \left(\text{blue circle} \text{---} \text{wavy line} \right) + \dots$$

Single-graviton production
amplitude (in-in)

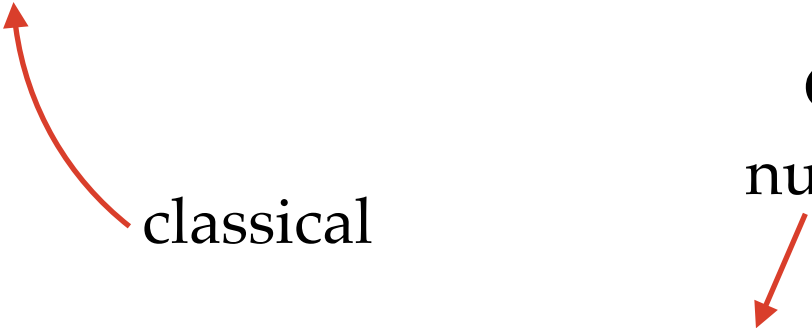
Amplitudes: quantum. Pair production? 

Motivation

Gravitational wave $\neq$ graviton?

classical

Graviton
number state


$$\langle n | \mathbb{h}_{\mu\nu} | n \rangle \sim \langle n | a(k) | n \rangle = 0$$

Resolution: wave is a coherent state of gravitons

$$S |\psi\rangle \supset \exp \left[\int_k \text{Single-graviton production amplitude} \right] |0\rangle$$

Motivation

Classically:

$$S |\psi\rangle \supset \exp \left[\int_k \text{blue circle} \text{ wavy line } a^\dagger(k) \right] |0\rangle$$

Coherent state
Classical waveform

Uncorrelated single graviton production

Semiclassically:

$$S |\psi\rangle \supset \exp \left[\int_{k_1, k_2} \text{blue circle} \text{ wavy line } a^\dagger(k_1) a^\dagger(k_2) \right] |0\rangle$$

Squeezed state

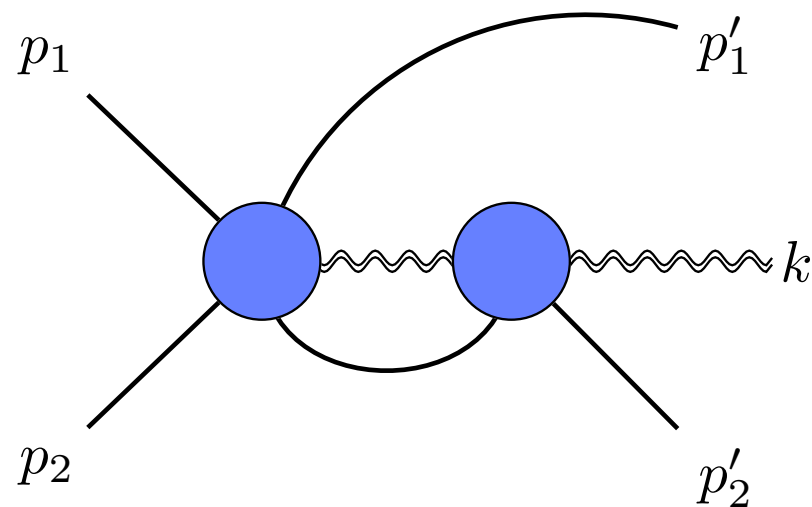
Uncorrelated production of pairs

Hawking radiation: pair production from BH

Motivation

So study how to compute Hawking radiation from amplitudes

Eg usual radiation backreaction

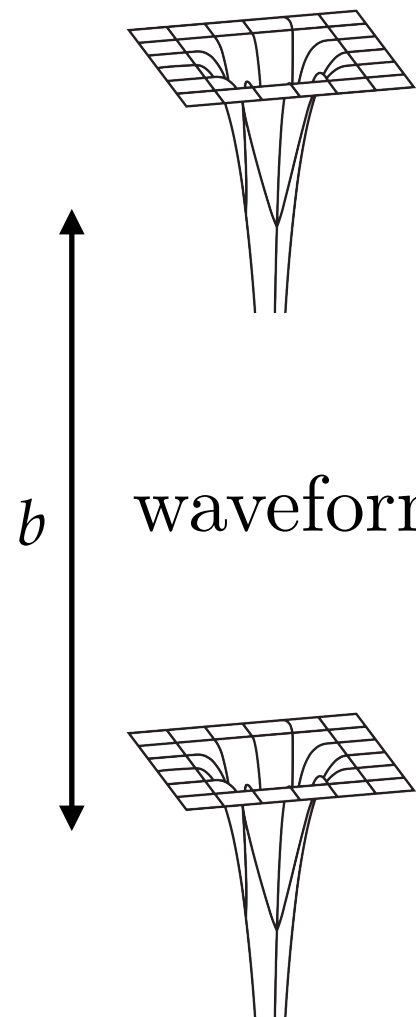


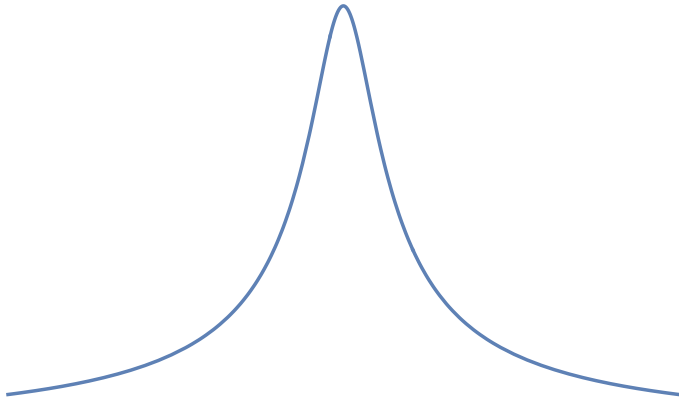
Compute reaction to Hawking pairs?

This talk

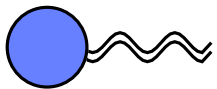
1. Waveforms and (generalised) amplitudes
2. Amplitudes and Bogoliubov
3. Hawking's scattering process, spectrum
4. Outlook

Waveforms & Amplitudes



$\text{waveform}(t) =$


$= \int \langle p'_1 p'_2 | S^\dagger a(k) S | p_1 p_2 \rangle + \dots$

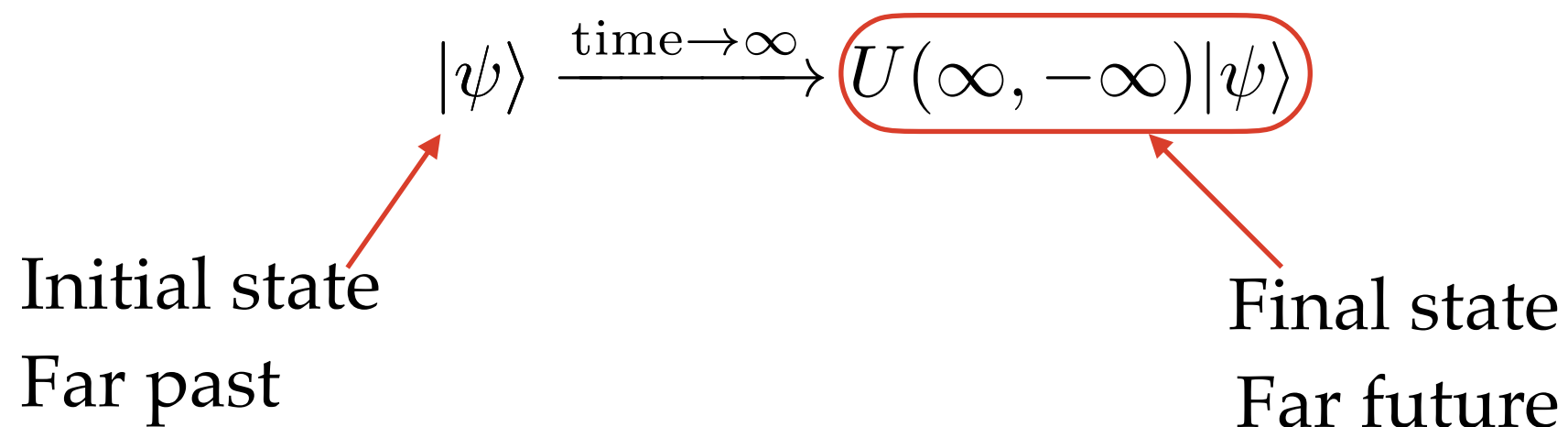
Generalised amplitude 

Classical point particle approximation (corrected at high order)

Kosower, Maybee & DOC
Cristofoli, Gonzo, Kosower & DOC

Waveforms & Amplitudes

Amplitudes: from time evolution (interaction picture)



Waveforms & Amplitudes

Amplitudes: from time evolution (interaction picture)

$$|\psi\rangle \xrightarrow{\text{time} \rightarrow \infty} S|\psi\rangle \quad U(-\infty, \infty) = S = 1 + iT$$

Matrix elements of T are the amplitudes

$$\langle q_1 \cdots q_m | T | p_1 \cdots p_n \rangle = \mathcal{A}(p_1 \cdots p_n \rightarrow q_1 \cdots q_m) \delta^4(\text{total momentum})$$

Waveforms & Amplitudes

Measure *expectation* of curvature component in classical limit

Riemann curvature
operator

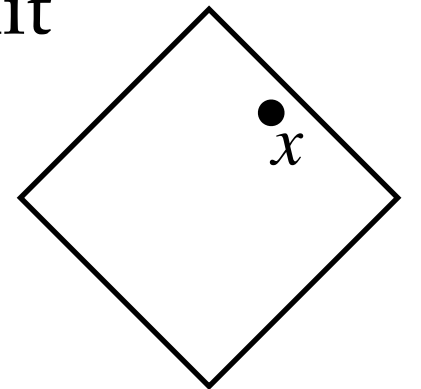
$$\text{waveform} \equiv \langle \psi | S^\dagger \mathbb{R} \dots (x) S | \psi \rangle$$

Final state
Amplitudes!

Dominant term at large distances
(Ψ_4)

$$|\psi\rangle \sim \int \mathcal{N} \mathcal{N} e^{ip_1 \cdot b} |p_1 p_2\rangle$$

Classical initial conditions



Waveforms & Amplitudes

Measure *expectation* of curvature component in classical limit

$$\mathbb{R}....(x) = \partial.\partial.\mathbb{h}..(x)$$

Leading in distance

Graviton polarisation

$$\text{waveform} = \int d\Phi(k) \left[k k \varepsilon \varepsilon e^{-ik \cdot x} \langle \psi | S^\dagger a(k) S | \psi \rangle + \text{c.c.} \right]$$

On-shell
phase space

Waveforms & Amplitudes

$$\text{waveform} = \int d\Phi(k) [kk \varepsilon \varepsilon e^{-ik \cdot x} \langle \psi | S^\dagger a(k) S | \psi \rangle + \text{c.c.}]$$

Key object to understand

$$\langle \psi | S^\dagger a(k) S | \psi \rangle = \int_{p_i, p'_i} \psi^*(p'_1, p'_2) \psi(p_1, p_2) \underbrace{\langle p'_1, p'_2 | S^\dagger a(k) S | p_1, p_2 \rangle}_{\text{Generalised amplitude}}$$

*Caron-Huot, Giroux, Hannesdóttir, Mizera
Elkhidir, Sergola, Vazquez-Holm, DOC*

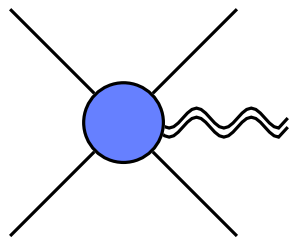
Waveforms & Amplitudes

Generalised amplitude closely related to amplitudes

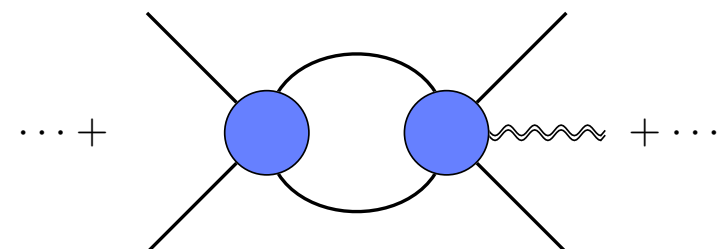
$$\mathcal{E} \delta^4(\text{momentum}) \equiv -i \langle p'_1 p'_2 | S^\dagger a(k) S | p_1 p_2 \rangle$$

$$= \langle p'_1 p'_2 k | T | p_1 p_2 \rangle - i \langle p'_1 p'_2 | T^\dagger a(k) T | p_1 p_2 \rangle$$

5 point amplitude



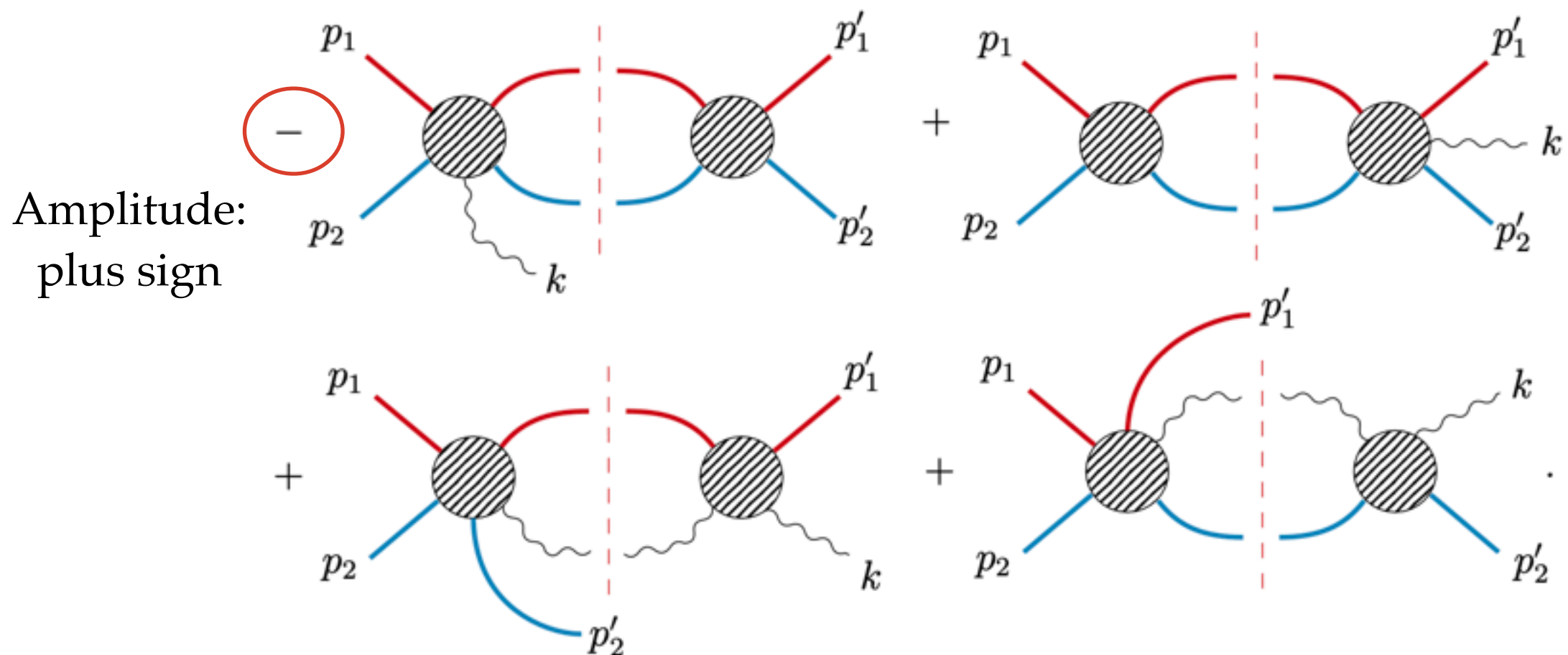
Cut (subtraction) term



Waveforms & Amplitudes

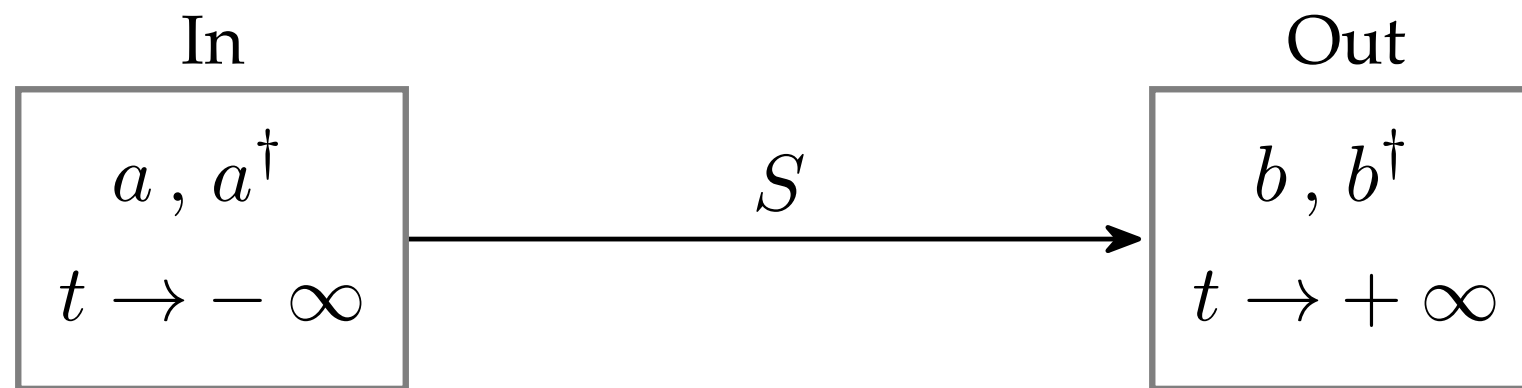
At NLO only difference:

$$2 \operatorname{Im}' \mathcal{E}(p_1 p_2 \rightarrow p'_1 p'_2 k_\eta) \hat{\delta}^D(p_1 + p_2 - p'_1 - p'_2 - k) =$$



Waveforms & Amplitudes

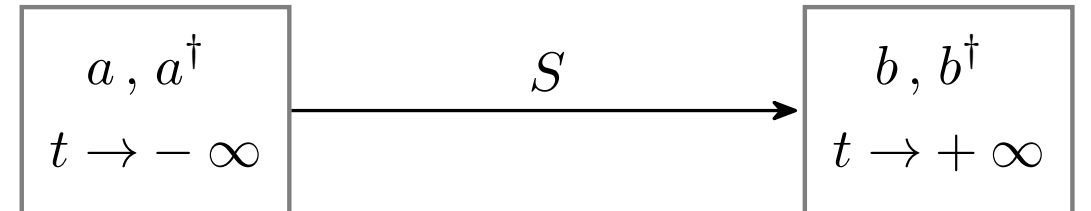
Past/future ladder operators



$$b(p) = S^\dagger a(p) S$$

Waveforms & Amplitudes

Past/future ladder operators



$$i\mathcal{A}\delta^4(\text{momentum}) = \langle 0|b(p'_1)b(p'_2)b(k)a^\dagger(p_1)a^\dagger(p_2)|0\rangle$$

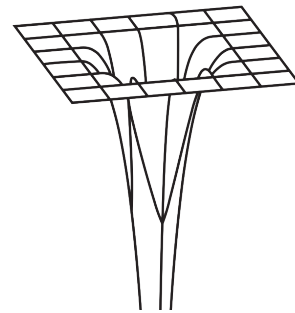
$$i\mathcal{E}\delta^4(\text{momentum}) = \langle 0|a(p'_1)a(p'_2)b(k)a^\dagger(p_1)a^\dagger(p_2)|0\rangle$$

Physics: different $i\epsilon$ prescription. In-out vs in-in, Schwinger-Keldysh

Expectation admits classical limit, amplitudes doesn't

Waveforms & Amplitudes

Scattering on background:



●
$$\text{waveform}(t) = \int \langle \Omega | a(p'_1) S^\dagger a(k) S a^\dagger(p_1) | \Omega \rangle$$

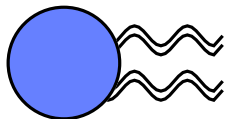
Generalised 3-point amplitude
on background

Appropriate at LO in m/M

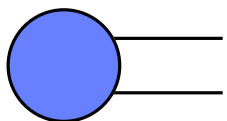
Amplitudes & Bogoliubov

Amplitudes & Bogoliubov

Pair production: any species



Graviton pair production



Scalar pair production: simpler

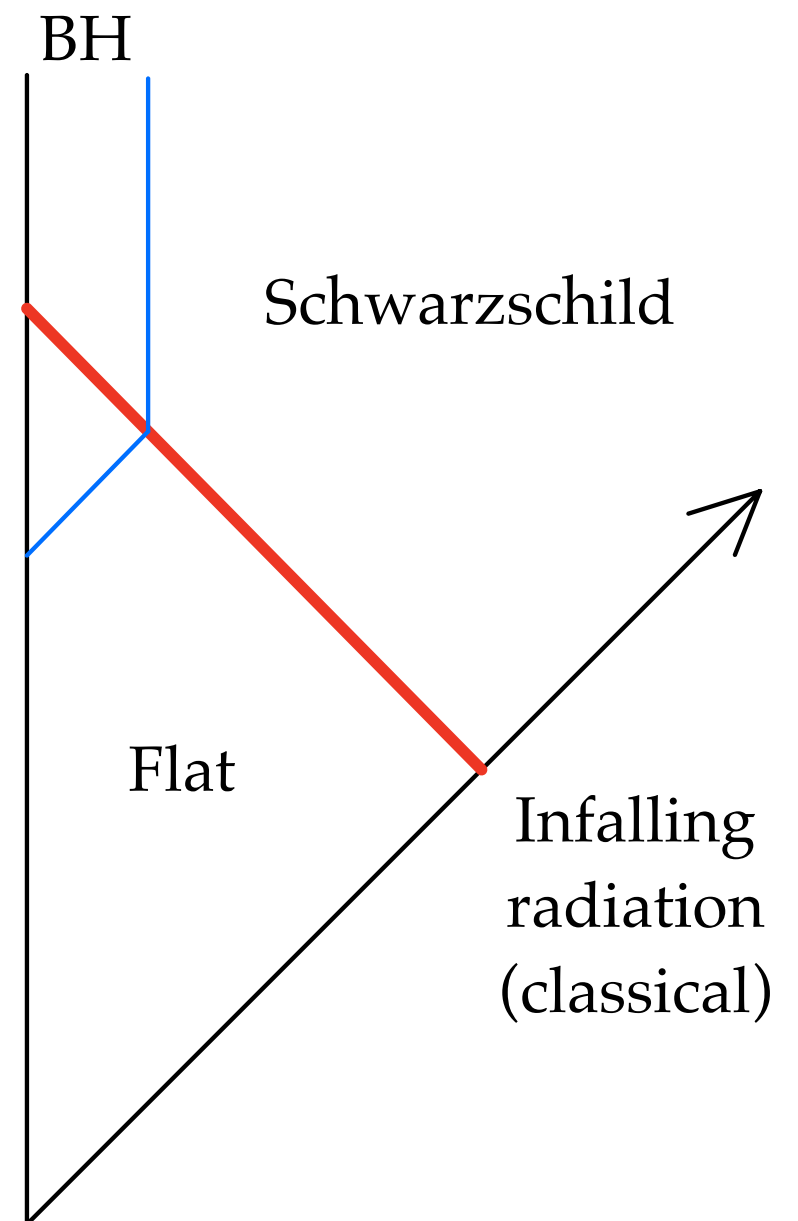
Amplitudes & Bogoliubov

Produced from dynamical background metric “Vaidya”

Thin shell of radiation
falling in to origin

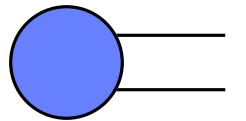
$$g^{\mu\nu} = \eta^{\mu\nu} + \frac{2GM\Theta(t+r)}{r} k^\mu k^\nu$$

Kerr-Schild vector
 $k_\mu dx^\mu = dt + dr$

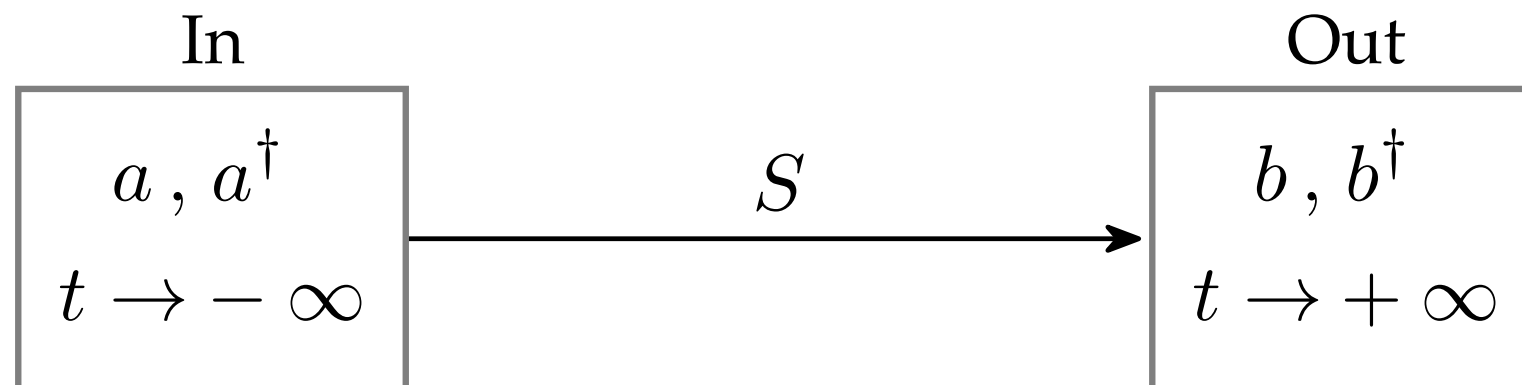


Amplitudes & Bogoliubov

$$\text{action} = \frac{1}{2} \int d^4x \partial_\mu \phi \partial_\nu \phi (\eta^{\mu\nu} + h^{\mu\nu})$$



Bogoliubov very useful:



$$b(p) = S^\dagger a(p) S$$

Amplitudes & Bogoliubov

Exact field eom $\partial^2 \phi + \partial_\mu h^{\mu\nu}(x) \partial_\nu \phi = 0$

Past solutions $P(x, p)$

$$P(x, p) \xrightarrow[t \rightarrow -\infty]{} e^{-ix \cdot p}$$

Future solutions $F(x, p)$

$$F(x, p) \xrightarrow[t \rightarrow \infty]{} e^{-ix \cdot p}$$

$$\phi(x) = \int d\Phi(p) [P(x, p)a(p) + \text{h.c.}]$$

$$\phi(x) = \int d\Phi(p) [F(x, p)b(p) + \text{h.c.}]$$

Amplitudes & Bogoliubov

$$b(k) = S^\dagger a(k) S = \int d\Phi(p) (A(k, p) a(p) + B(k, p) a^\dagger(p))$$



$$P \sim \Sigma (AF + \bar{B}\bar{F})$$

F is a basis



$$\phi(x) = \int d\Phi(p) [P(x, p) a(p) + \text{h.c.}] = \phi(x) = \int d\Phi(p) [F(x, p) b(p) + \text{h.c.}]$$

Amplitudes & Bogoliubov

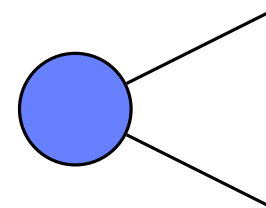
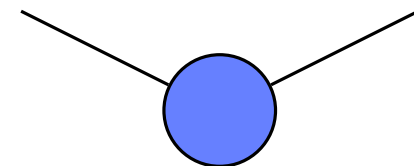
$$b(k) = S^\dagger a(k) S = \int d\Phi(p) (A(k, p) a(p) + B(k, p) a^\dagger(p))$$

Bogoliubov coefficients = (generalised) amplitudes

Vaidya vacuum

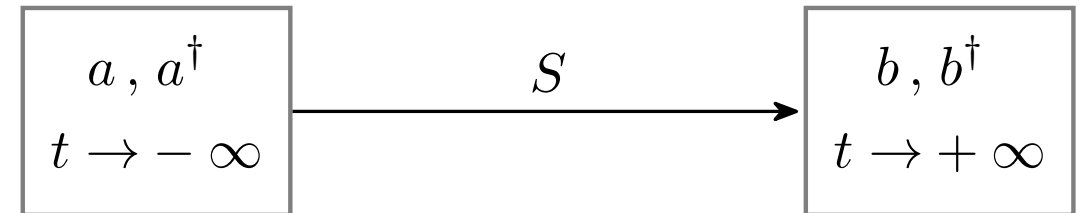

$$A(k, p) = \langle \Omega | S^\dagger a(k) S a^\dagger(p) | \Omega \rangle$$

$$B(k, p) = \langle \Omega | a(p) S^\dagger a(k) S | \Omega \rangle$$



Amplitudes & Bogoliubov

Count number of particles in future



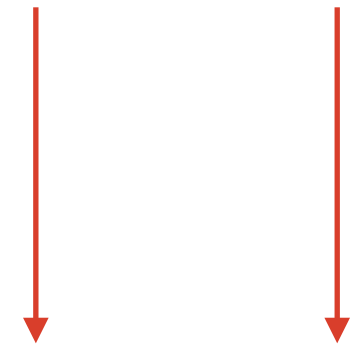
$$\begin{aligned}\langle \Omega | b^\dagger(p) b(p) | \Omega \rangle &= \langle \Omega | S^\dagger a^\dagger(p) a(p) S | \Omega \rangle \\ &= \int d\Phi(k) \bar{B}(p, k) B(p, k)\end{aligned}$$

Just need Bogoliubov B

Amplitudes & Bogoliubov

“Crossing”: follows since F and P are both bases $P \sim \Sigma (AF + \bar{B}\bar{F})$

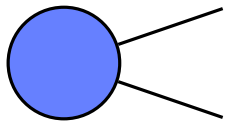
$$A(k, p) = \int d^3x \left[F^*(x, k) i\partial_t e^{\ominus ip \cdot x} - e^{\ominus ip \cdot x} i\partial_t F^*(x, k) \right]$$

$$B(k, p) = \int d^3x \left[F^*(x, k) i\partial_t e^{\oplus ip \cdot x} - e^{\oplus ip \cdot x} i\partial_t F^*(x, k) \right]$$


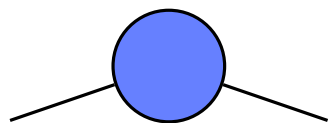
Analytically continue $p \rightarrow -p$

Amplitudes & Bogoliubov

“Crossing”



Scalar pair production



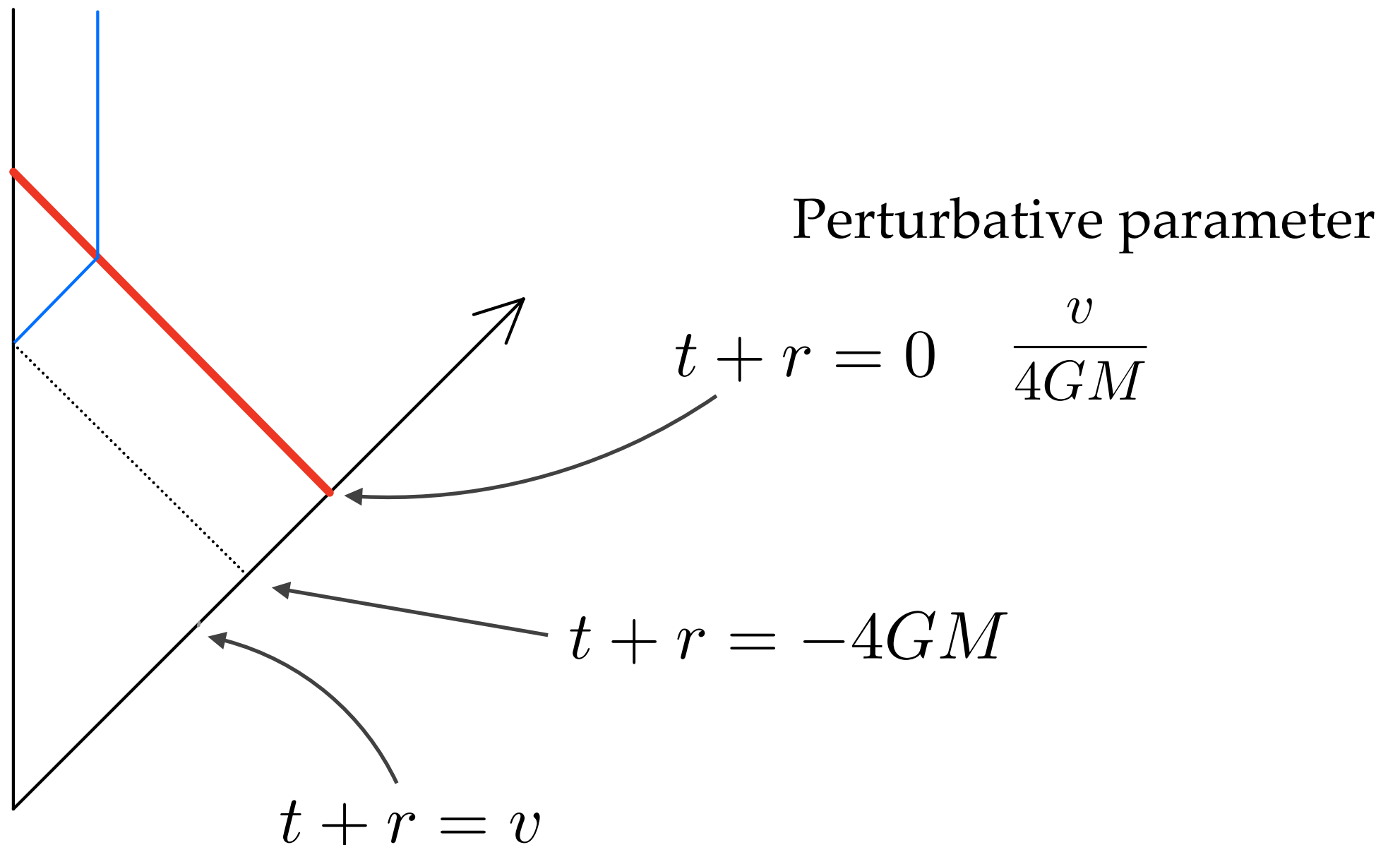
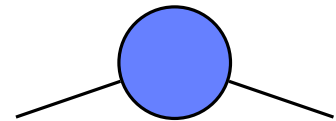
Scalar scattering: more familiar

“Hawking” scattering
Quantum process

Hawking's Scattering Process

Hawking Scattering

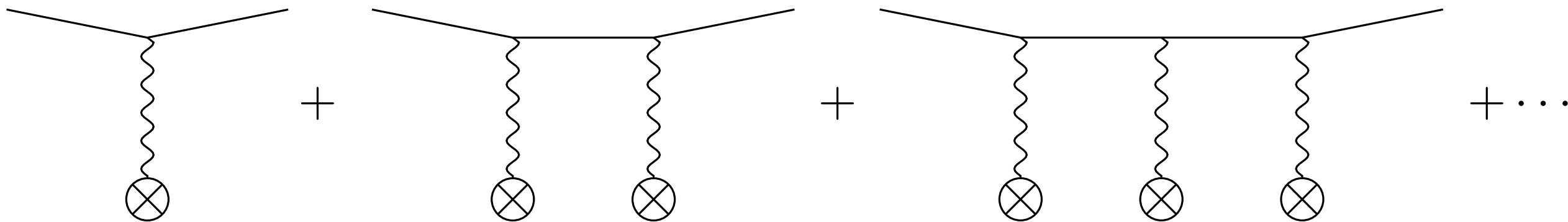
Scatter spherical particle through background



Hawking Scattering

Lippmann-Schwinger approach

$$\partial^2 \phi + \partial_\mu h^{\mu\nu}(x) \partial_\nu \phi = 0$$

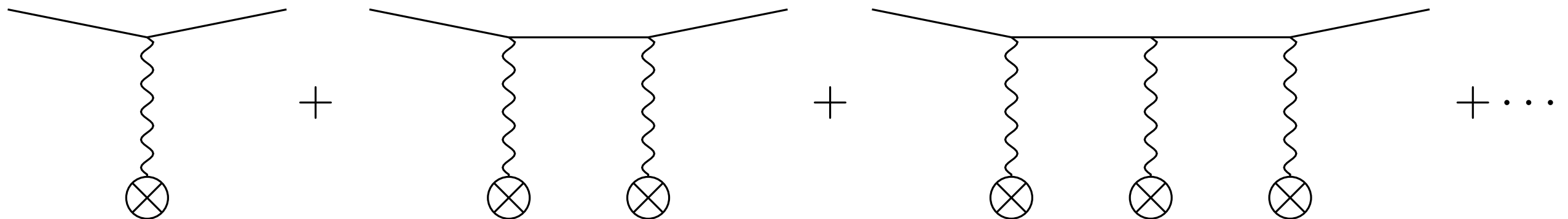


Hawking Scattering

Exact quantum scattering: hard

Simplify: high energy “geometric optics” $\sim e^{iS(x)/\eta}$

Analogous to $\hbar$

$$\mathcal{A}_0^{(-1)} \sim \frac{1}{\eta} \quad \mathcal{A}_1^{(-2)} \sim \frac{1}{\eta^2} \quad \mathcal{A}_2^{(-3)} \sim \frac{1}{\eta^3} \quad + \dots$$


The diagram illustrates the expansion of Hawking scattering amplitudes. It shows three terms in a series, each representing a different order of interaction. The first term, $\mathcal{A}_0^{(-1)} \sim \frac{1}{\eta}$, shows a single wavy line connecting a vertex to a circle with an 'X'. The second term, $\mathcal{A}_1^{(-2)} \sim \frac{1}{\eta^2}$, shows two wavy lines connecting a vertex to two such circles. The third term, $\mathcal{A}_2^{(-3)} \sim \frac{1}{\eta^3}$, shows three wavy lines connecting a vertex to three such circles. The series continues with an ellipsis. The vertices are represented by lines that split or join at a point.

Hawking Scattering

Scatter massless spherical particle through background

$$|\psi\rangle = \int d\Phi(p) \varphi(p) |p\rangle = \int d\Phi(p) \int dv \varphi(v) e^{iEv} |p\rangle$$

$$= \int d\Phi(p) \int dv \varphi(v) e^{ip \cdot b(v)} |p\rangle$$

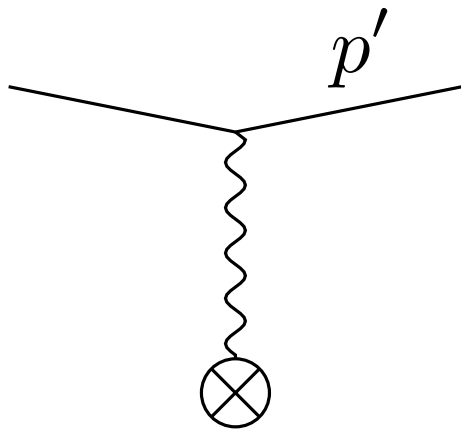
Time-like “impact
parameter”

Spherical symmetry

$$E = |\mathbf{p}|$$

Hawking Scattering

Tree



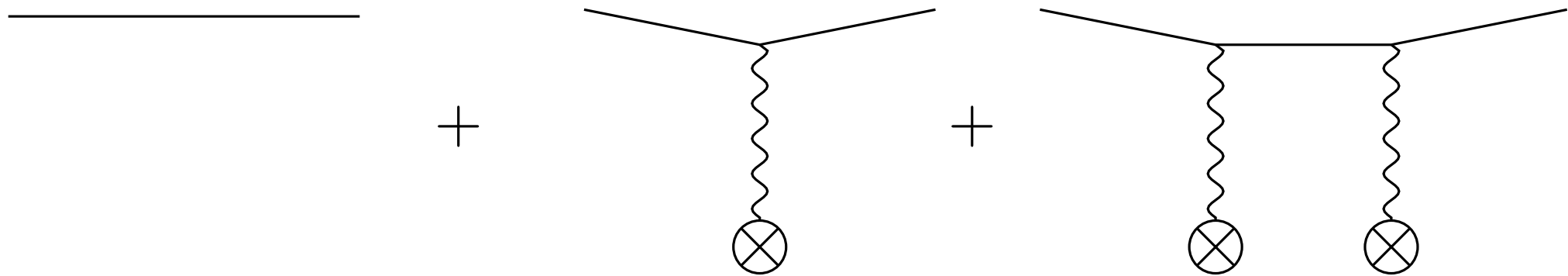
$$S_{\text{int}} = \frac{1}{2} \int d^4x h^{\mu\nu} \partial_\mu \phi \partial_\nu \phi$$

Compute using GW methods:

$$\text{diagram} = \int dv \varphi(v) e^{ip' \cdot b(v)} \left(-4GM E' \log(-v/\mu) \right)$$

Hawking Scattering

Iterates & exponentiates



$$\mathcal{A} = \mathcal{N} \int dv \phi(v) e^{i(E'p' - E_0)v} \exp\left(-\frac{4iGM E' \log(-v/\mu)}{4iGM E' \log(-v/\mu)}\right)$$

Spherical energy eigenstate $|\psi\rangle = \frac{1}{4\pi} \int d\Omega |E_0, E_0 \mathbf{n}\rangle$

Hawking Scattering

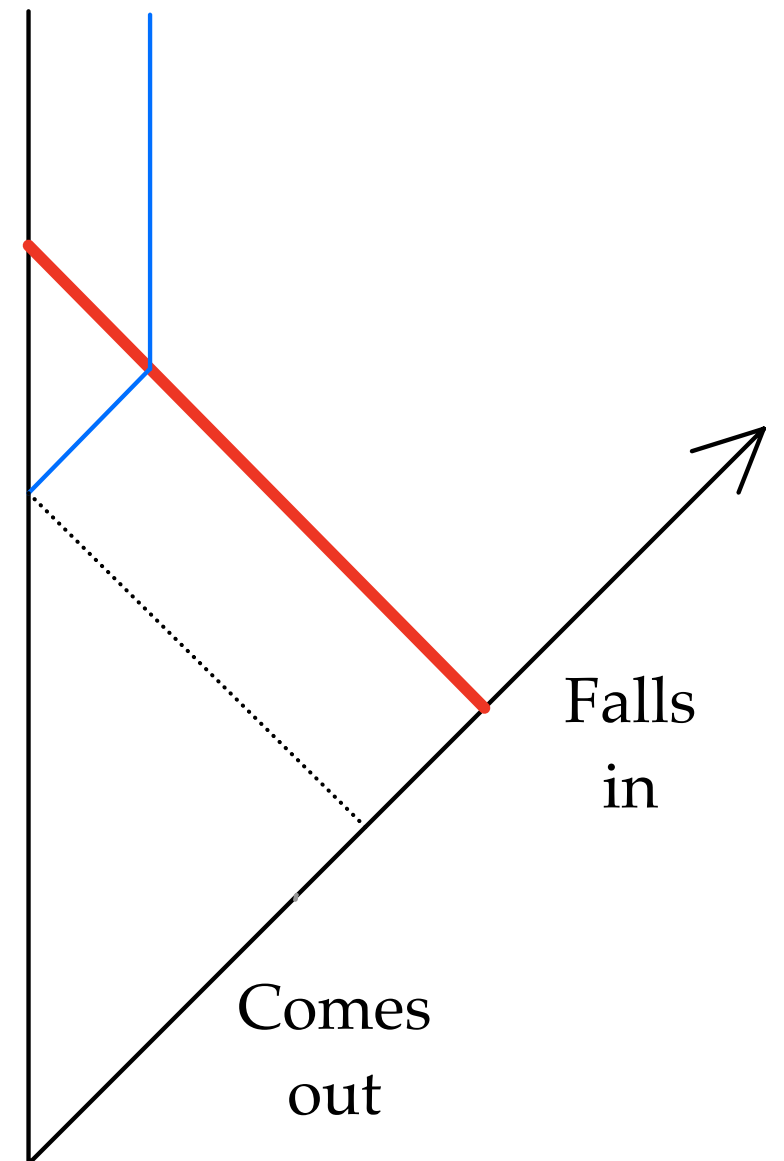
Not fully inclusive

$$\mathcal{A} = \mathcal{N} \int dv e^{i(E' - E_0)v} e^{-4iGM E' \log(-v/\mu)}$$

Horizon

$$\mathcal{A} = \mathcal{N} \int_{-\infty}^0 dv e^{i(E' - E_0)v} e^{-4iGM E' \log(-v/\mu)}$$

Hawking's amplitude



Pair Production

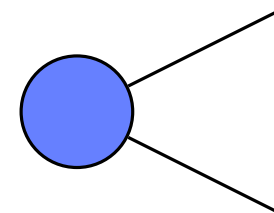
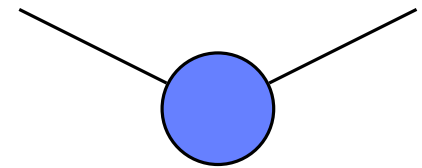
Simply cross in-state to out-state

$$\mathcal{A} = \mathcal{N} \int_{-\infty}^0 dv e^{i(E' - E_0)v} e^{-4iGM E' \log(-v/\mu)}$$

Crossing

$$\mathcal{B} = \mathcal{N} \int_{-\infty}^0 dv e^{i(E' + E_0)v} e^{-4iGM E' \log(-v/\mu)}$$

$$\sim e^{-2\pi GME'} \Gamma(1 - 4iGME')$$



Pair Production

Count outgoing states

$$|\mathcal{B}|^2 = (\text{factor}) \frac{1}{e^{8\pi G M E'} - 1}$$



Thermal distribution

$$T = \frac{1}{8\pi G M}$$

Outlook

Higher Orders

Analog of PM expansion

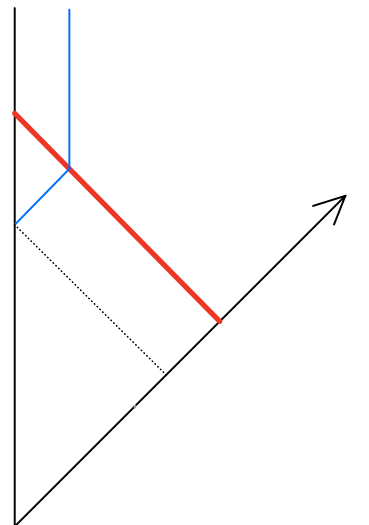
$$\exp\left(i\frac{S}{\eta}\right) = \exp\left(i\frac{S^{(0)} + S^{(1)} + \dots}{\eta}\right)$$

Also computed $S^{(1)}$

$$S^{(0)} + S^{(1)} = -4GM E' \log(-v/\mu) - \frac{16G^2 M^2 E'}{v}$$

$$\simeq -4GM E' \log(-(v + 4GM)/\mu)$$

Finite size effect



Finite radius of horizon

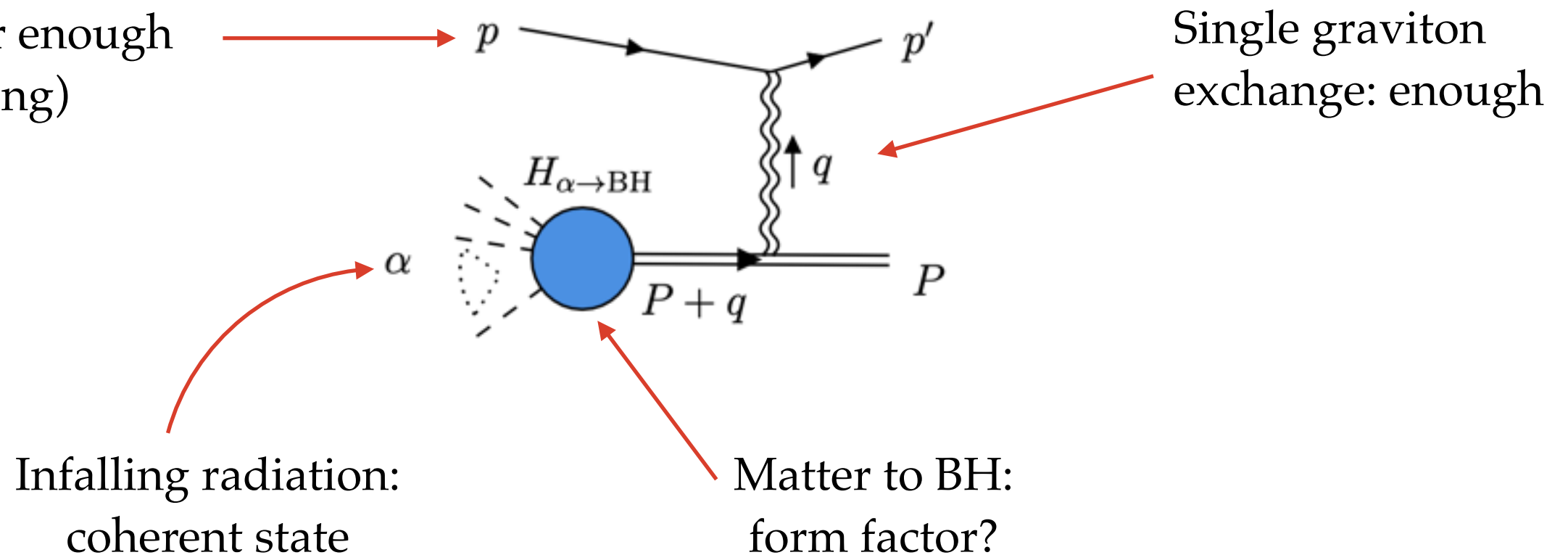
Outlook

Hawking computation: closely related to PM methods of amplitudes

❖ Remove background?

Pair production: crossing

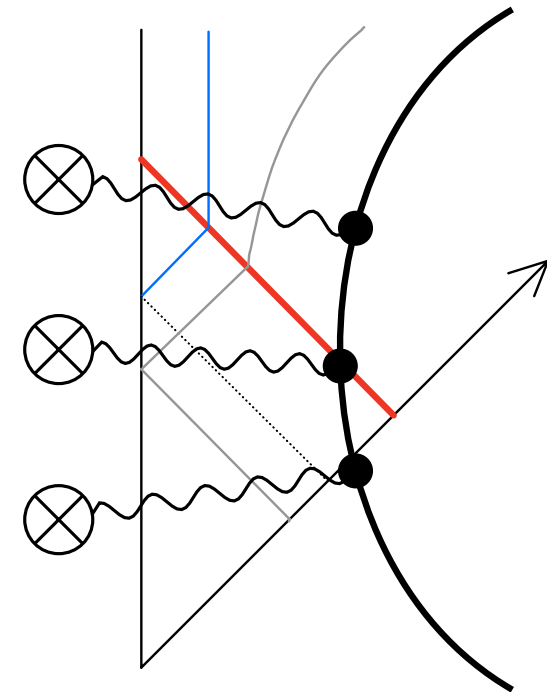
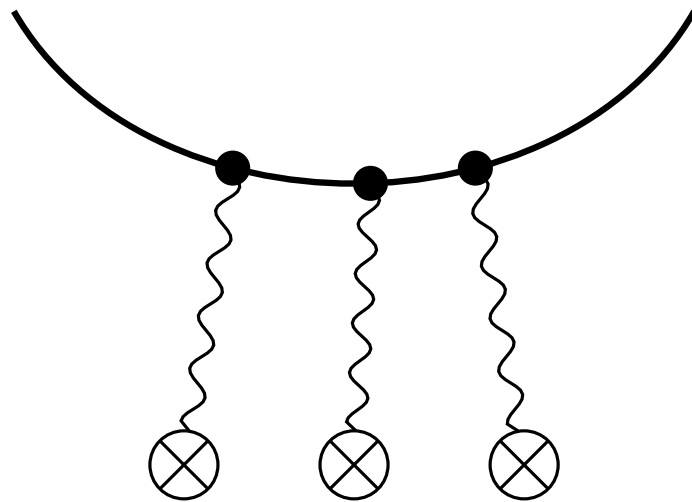
One pair enough
(Squeezing)



Outlook

Hawking computation: closely related to PM methods of amplitudes

- ❖ Remove background?
- ❖ Non-local pair production?



Outlook

Hawking computation: closely related to PM methods of amplitudes

- ❖ Remove background?
- ❖ Non-local pair production - tunneling?
- ❖ Exponentiate with more care: stationary phase
- ❖ Radiation reaction?
- ❖ Page curve?

Parikh & Wilczek