High-Energy Behavior of Scattering Amplitudes in Theories with Purely Virtual Particles

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- **1** Introduction on PVP
- **2** Scattering in quantum gravity
- **3** Unitarity Bounds and Perturbativity
- **4** O(N) models and large-N expansion

#### **5** Results

# Introduction on PVP

# Particle Physics and Quantum Gravity

In particle physics

Locality, renormalizability and unitarity (plus symmetries)

 $\Downarrow$ 

Standard Model

# Particle Physics and Quantum Gravity

#### In particle physics

#### Locality, renormalizability and unitarity (plus symmetries)

#### $\Downarrow$

#### Standard Model

In gravity

Einstein gravity  $S = \int \sqrt{-g}R$ Unitary  $\checkmark$  Renormalizable  $\times$  Stelle gravity  $S = \int \sqrt{-g} \left[ R + R^2 + R^2_{\mu\nu} \right]$ Unitary × Renormalizable ✓ Degrees of freedom in Stelle theory

$$S_{\rm QG}(g) = -\frac{1}{16\pi} \int \sqrt{-g} \left( M_{\rm pl}^2 R - \frac{\xi}{6} R^2 + \frac{\alpha}{2} C_{\mu\nu\rho\sigma} C^{\mu\nu\rho\sigma} \right), \qquad g_{\mu\nu} = \eta_{\mu\nu} + \sqrt{32\pi} h_{\mu\nu}.$$
DOF

Massless graviton Massive scalar  $\phi$  Massive spin-2 ghost  $\chi$ 

 $h_{\mu\nu}$  propagator

$$\begin{split} \frac{i}{p^2} \left[ \frac{\Pi_{\mu\nu\rho\sigma}^{(2)}}{(M_{\rm pl}^2 - \alpha p^2)} - \frac{\Pi_{\mu\nu\rho\sigma}^{(0)}}{2(M_{\rm pl}^2 - \xi p^2)} \right] + \text{g.f.} &= \frac{i}{2M_{\rm pl}^2} \left[ \frac{2\Pi_{\mu\nu\rho\sigma}^{(2)} - \Pi_{\mu\nu\rho\sigma}^{(0)}}{p^2} - \frac{2\Pi_{\mu\nu\rho\sigma}^{(2)}}{p^2 - m_\chi^2} + \frac{\Pi_{\mu\nu\rho\sigma}^{(0)}}{p^2 - m_\chi^2} \right] + \text{g.f.}, \\ m_\chi^2 &= M_{\rm pl}^2 / \alpha, \qquad m_\phi^2 = M_{\rm pl}^2 / \xi. \end{split}$$

Off-shell  $\chi$  and  $\phi$  are responsible for renormalizability  $\checkmark$ 

On-shell  $\chi$  is responsible for the violation of unitarity X

# **Purely Virtual Particles**

D. Anselmi and MP, JHEP 06 (2017) 066, D. Anselmi JHEP 06 (2017) 086

In quantum field theory a particle can be either "virtual" or "real"(on shell)

Real photon: light

Virtual photon: electromagnetic force



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On-shell and virtual parts are related because of the optical theorem and the Feynman prescription



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#### Purely virtual particles (PVP)

- Only mediate interactions
- Contribute to renormalization
  - $\bullet$  Never be on shell
  - Consistent with unitarity.

# • Optical theorem

$$S = 1 + iT$$
,  $SS^{\dagger} = 1$   $\Leftrightarrow$   $-i(T - T^{\dagger}) = TT^{\dagger}$ .

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• Goal: consistent projection

$$\int d\Pi_f \left| \underbrace{-F}_{away} \right|^2 = 2 \operatorname{Im} \left[ (-i) \underbrace{-F}_{away} \right] = \underbrace{-F}_{away} = 0 \text{ by fakeon prescription}$$

# Cuts and imaginary parts

PVP cut propagator vanishes. A bubble diagram has only one threshold



More complicated diagrams have multiple thresholds  $\Rightarrow$  Modified functions

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More complicated diagrams have multiple thresholds  $\Rightarrow$  Modified functions



 $Diagram \to Diagram - \Delta_{Diagram}^{n}, \qquad n = \# of PVP \text{ inside the loop}$ (1)

For one-loop bubble diagram (this talk):  $\Delta_{\text{Bubble}}^{1,2} = \text{Re(Bubble})$ For triangles and boxes see A. Melis and MP, PRD 108 (2023) 9, 096021. For general procedure D. Anselmi JHEP 11 (2021) 030.

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# Threshold decomposition and spectral identities

D. Anselmi, JHEP 11 (2021) 030

$$G_N = \int \frac{\mathrm{d}^D k}{(2\pi)^D} \prod_{a=1}^N \frac{1}{(k-p_a)^2 - m_a^2 + i\epsilon_a} = \int \frac{\mathrm{d}^{D-1} \mathbf{k}}{(2\pi)^{D-1}} \left(\prod_{a=1}^N \frac{1}{2\omega_a}\right) G_N^s$$

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Skeleton diagram

$$G_N^s = \int \frac{\mathrm{d}k^0}{2\pi} \prod_{a=1}^N \frac{2\omega_a}{(k^0 - e_a)^2 - \omega_a^2 + i\epsilon_a}, \qquad e_a = p_a^0, \qquad \omega_a = \sqrt{(\mathbf{k} - \mathbf{p}_a)^2 + m_a^2}$$

Spectral identities  $G^{s} + (G^{s})^{*} + \sum_{cuts} G_{c}^{s} = 0.$ which holds threshold by thresholds.

$$B^{s} = -\frac{i}{e_{1} - e_{2} - \omega_{1} - \omega_{2} + i\epsilon} - \frac{i}{e_{2} - e_{1} - \omega_{1} - \omega_{2} + i\epsilon}.$$

Using  $\frac{i}{x+i\epsilon} = \mathcal{P}\frac{i}{x} + \pi\delta(x)$  everywhere

$$B^s = -i\mathcal{P}_2 - \Delta^{12} - \Delta^{21}$$

$$\mathcal{P}^{ab} = \mathcal{P}\frac{1}{e_a - e_b - \omega_a - \omega_b}, \qquad \mathcal{P}_2 = \mathcal{P}^{ab} + \mathcal{P}^{ba}, \qquad \Delta^{ab} = \pi\delta(e_a - e_b - \omega_a - \omega_b)$$

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Diag. Terms	$\bigcirc^2$			
$\mathcal{P}_2$	-i	i	0	0
$\Delta^{12}$	-1	-1	2	0
$\Delta^{21}$	-1	-1	0	2

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$$\boxed{\begin{array}{c|c} \text{Diag.} & 2 \\ \text{Terms} & 1 \\ \hline \mathcal{P}_2 & -i \\ \hline \mathcal{Q}_1 & -1 \\ \hline \mathcal{Q}_2 & -1$$

If we want to include PVP, we kill all the  $\Delta$ 's that contain at least one PVP frequency.

Diag. Terms	2 3 1		X.	Tundu				
$\mathcal{P}_3$	-i	i	0	0	0	0	0	0
$\Delta^{12} Q^{13}$	-1	-1	2	0	0	0	0	0
$\Delta^{23} Q^{21}$	$^{-1}$	-1	0	2	0	0	0	0
$\Delta^{31} Q^{32}$	$^{-1}$	-1	0	0	2	0	0	0
$\Delta^{21} Q^{23}$	-1	-1	0	0	0	2	0	0
$\Delta^{32} Q^{31}$	-1	-1	0	0	0	0	2	0
$\Delta^{13} Q^{12}$	-1	-1	0	0	0	0	0	2
$\Delta^{12}\Delta^{13}$	i	-i	2i	0	0	0	0	-2i
$\Delta^{23}\Delta^{21}$	i	-i	0	2i	0	-2i	0	0
$\Delta^{31}\Delta^{32}$	i	-i	0	0	2i	0	-2i	0
$\Delta^{21}\Delta^{31}$	i	-i	0	0	2i	-2i	0	0
$\Delta^{32} \Delta^{12}$	i	-i	2i	0	0	0	-2i	0
$\Delta^{13} \overline{\Delta^{23}}$	i	-i	0	2i	0	0	0	-2i

$$\mathcal{P}_3 = \mathcal{P}^{12} \mathcal{P}^{13} + \text{cycl} + (e \to -e), \qquad \mathcal{Q}^{ab} = \mathcal{P}^{ab} - \mathcal{P} \frac{1}{e_a - e_b - \omega_a + \omega_b}.$$

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$$\mathcal{P}_{ABC} = \mathcal{P}^{12}\mathcal{P}^{13} + \text{cycl} + (e \to -e), \qquad \mathcal{Q}^{ab} = \mathcal{P}^{ab} - \mathcal{P}\frac{1}{e_a - e_b - \omega_a + \omega_b}.$$

# Scattering in quantum gravity

# Graviton scattering in quantum gravity

In Einstein gravity (unitary, nonrenormalizable)

$$S_{\rm H}(g) = -\frac{1}{2\kappa^2} \int \sqrt{-g}R, \qquad \kappa^2 = 8\pi G, \qquad g_{\mu\nu} = \eta_{\mu\nu} + 2\kappa h_{\mu\nu} \tag{2}$$

Tree level amplitudes

$$\mathcal{A}_{hh \to hh}^{\mathrm{EG}} \sim \kappa^2 s, \qquad s = \text{ c.o.m. energy squared}$$
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In Stelle gravity (non-unitary, renormalizable)

$$S_{\rm SG}(g) = -\frac{1}{2\kappa^2} \int \sqrt{-g} \left[ \zeta R - \frac{1}{6\xi} R^2 + \frac{1}{2\alpha} C_{\mu\nu\rho\sigma} C^{\mu\nu\rho\sigma} \right], \qquad C^{\mu}{}_{\nu\rho\sigma} = \text{Weyl tensor} \tag{4}$$

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P. Donà, S. Giaccari, L. Modesto, L. Rachwal, Y. Zhu JHEP 08 (2015) 038.

$$\mathcal{A}_{hh\to hh}^{\rm SG} = \mathcal{A}_{hh\to hh}^{\rm EG} \tag{5}$$

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# Equivalent formulation of Stelle gravity

D. Anselmi and MP, JHEP 11 (2018) 021.

$$S_{\rm SG}(g) = -\frac{1}{2\kappa^2} \int \sqrt{-g} \left[ \zeta R - \frac{1}{6\xi} R^2 + \frac{1}{2\alpha} C_{\mu\nu\rho\sigma} C^{\mu\nu\rho\sigma} \right] \tag{6}$$

Equivalent action

auxiliary fields  $\phi$ ,  $\chi_{\mu\nu}$  + Weyl transformation + field redefinitions.  $\downarrow$ 

$$S(g,\phi,\chi) = -\frac{1}{2\kappa^2} \int \sqrt{-g} \left[ \zeta R + \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right] + S_{\chi\chi}(g,\chi) + S_{\rm int}(g,\chi,\phi)$$
(7)  
$$V(\phi) = \frac{m_\phi^2}{\kappa^2} (1 - e^{\tilde{\kappa}\phi})^2, \qquad S_{\chi\chi} = -S_{\rm Pauli-Fierz}$$

Vertices in  $S_{int}$ 

# Unitarity Bounds and Perturbativity

# Unitarity bounds and perturbativity

Usual derivation of Unitarity bounds for 2-to-2 scattering of scalars

$$\sigma(s) = \frac{1}{32\pi s} \int_{-1}^{1} \mathrm{d}v |\mathcal{M}(s,v)|^2, \qquad v \equiv \cos\theta.$$
(10)

$$S = 1 + iT, \qquad S^{\dagger}S = 1, \qquad -i\left(T - T^{\dagger}\right) = \frac{1}{2}T^{\dagger}T \qquad (11)$$

$$\Downarrow$$

$$\operatorname{Im}\mathcal{M}(s,1) = \sqrt{\kappa(s,m_1^2,m_2^2)} \sum_X \sigma_X(s) \ge \sqrt{\kappa(s,m_1^2,m_2^2)} \sigma(s), \tag{12}$$

$$\kappa(x, y, z) = x^{2} + y^{2} + z^{2} - xy - xz - yz$$
(13)

Partial wave expansion in Legendre polynomials

$$\mathcal{M}(s,v) = 16\pi \sum_{j=0}^{\infty} (2j+1)A_j(s)P_j(v), \qquad \sigma(s) = \frac{16\pi}{s} \sum_{j=0}^{\infty} (2j+1)|A_j(s)|^2$$
(14)

# Unitarity bounds and perturbativity

$$\sum_{j=0}^{\infty} (2j+1) \operatorname{Im} A_j(s) \ge \frac{\sqrt{\kappa(s, m_1^2, m_2^2)}}{s} \sum_{j=0}^{\infty} (2j+1) |A_j(s)|^2.$$
(15)

For elastic scattering only we have the equality, which leads to

$$|\mathcal{A}_j| \le 1, \qquad 0 \le \mathrm{Im}\mathcal{A}_j \le 1, \qquad |\mathrm{Re}\mathcal{A}_j| \le \frac{1}{2}, \qquad \mathcal{A}_j(s) = \frac{\sqrt{\kappa(s, m_1^2, m_2^2)}}{s} A_j(s) \tag{16}$$

So far Unitarity  $\Rightarrow$  (15). Therefore, if (15) is violated  $\Rightarrow$  Unitarity is violated.

# Unitarity bounds and perturbativity

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• Typical argument:

 $\mathcal{M} = \mathcal{M}^{\text{tree}} + \mathcal{M}^{\text{loops}}, \quad \mathcal{M}^{\text{tree}} = 16\pi A_0^{\text{tree}} + \dots, \quad A_0^{\text{tree}} \sim s \text{-channel diagrams}$ (17)

X if  $\mathcal{A}_0^{\text{tree}}$  violates (16) at some scale, then the theory violates unitarity at some scale.

Unitarity **and** perturbativity  $(|\mathcal{M}^{\text{loop}}| < |\mathcal{M}^{\text{tree}}|) \Rightarrow (16)$  for  $\mathcal{A}_0^{\text{tree}}$ 

 $\checkmark$  if  $\mathcal{A}_0^{\text{tree}}$  violates (16) at some scale, then **either** unitarity or perturbativity is violated at some scale

## Diagrammatic optical theorem

Unitarity can be checked by means of the cutting equations

$$2\mathrm{Im}\left(-iG\right) = -\sum_{c} G_{c} \tag{18}$$

which hold for any diagram G for any local QFT.

- Absence of ghosts  $\Rightarrow$  (18) is the diagrammatic version of the unitarity equation.
- The cutting equations (18) relies only on the diagramamtic expansion.
- No assumptions on the behavior of G.

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 $\Downarrow$ 

If  $\mathcal{A}_0^{\text{tree}}$  violates  $|\mathcal{A}_0^{\text{tree}}| \leq 1$  etc... at some scale, then perturbativity is violated at some scale.

# O(N) models

O(N) models and large-N diagrammatics

$$\mathcal{L}(\varphi) = \frac{1}{2} \partial_{\mu} \varphi^a \partial^{\mu} \varphi^a - \frac{1}{2} m^2 \varphi^a \varphi^a - \frac{g}{8} (\varphi^a \varphi^a)^2, \qquad a = 1, \dots N$$
(19)

Identify the orders in N

$$-\frac{g}{8}(\varphi^a\varphi^a)^2 \to \frac{1}{2g}\Omega^2 - \frac{1}{2}\Omega\varphi^a\varphi^a, \qquad \tilde{g} = gN \quad (\text{'t Hooft coupling}) \tag{20}$$

Study the limit  $N \to \infty$  by keeping  $\tilde{g}$  fixed.

- Each  $\Omega$  internal line gives 1/N
- Each closed  $\varphi$  loop gives N

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- Each closed  $\varphi$  loop gives N

Consider the scattering of  $\varphi^a\varphi^a\to\varphi^b\varphi^b$  with  $a\neq b$ 



All the bubble insertions are order 1/N and can be resummed

# Higher-derivative O(N) models

$$\mathcal{L}(\varphi) = \frac{1}{2} \partial_{\mu} \varphi^{a} F_{n} \left(\frac{-\Box}{M^{2}}\right) \partial^{\mu} \varphi^{a} - \frac{1}{2} m^{2} \varphi^{a} F_{n} \left(\frac{-\Box}{M^{2}}\right) \varphi^{a} - \frac{1}{8} \varphi^{2} G_{r} \left(\frac{-\Box}{M^{2}}\right) \varphi^{2}, \qquad \varphi^{2} \equiv \varphi^{a} \varphi^{a}, \quad (21)$$

$$F_n(z) = 1 + \sum_{i=1}^n f_i z^i, \qquad G_r(z) = \sum_{i=0}^r \lambda_i z^i,$$
 (22)

• Auxiliary fields and 't Hooft couplings

$$-\frac{1}{8}\varphi^2 G_r\left(\frac{-\Box}{M^2}\right)\varphi^2 \to \frac{1}{2}\Omega G_r^{-1}\left(\frac{-\Box}{M^2}\right)\Omega - \frac{1}{2}\Omega\varphi^2, \qquad \tilde{\lambda}_i = \lambda_i N$$
(23)

• Propagators

$$iD_{\Omega}(p^2) = iG_r(p^2), \qquad iD_{\varphi}^{ab}(p) = \frac{i\delta^{ab}}{(p^2 - m^2 + i\epsilon)F_n(p^2/M^2)} \equiv i\delta^{ab}D_{\rm HD}(p^2)$$
 (24)

# Bubble diagrams

$$D_{\rm HD}(p^2) = \frac{a_0}{p^2 - m^2 + i\epsilon} + \sum_{i=1}^n \frac{a_i}{p^2 - M_i^2 + i\epsilon}, \quad \text{with} \quad \sum_{i=0}^n a_i = 0$$
(25)

• Generic bubble diagram

$$B_{ij}(p^2) \equiv \int \frac{\mathrm{d}^D q}{(2\pi)^D} \frac{1}{(p+q)^2 - M_i^2 + i\epsilon} \frac{1}{q^2 - M_j^2 + i\epsilon},\tag{26}$$

$$B_{ij}(p^2,\tau) = \tau \operatorname{Re}B_{ij}(p^2) + i \operatorname{Im}B_{ij}(p^2), \qquad (\tau = 0 \text{ for PVP})$$
(27)

•  $\Omega$  self energy is  $N\Sigma(p^2, \tau)$ 

$$\Sigma(p^2,\tau) = \int \frac{\mathrm{d}^D q}{(2\pi)^D} D_{\mathrm{HD}}(p+q) D_{\mathrm{HD}}(q) = \sum_{i,j=0}^n a_i a_j B_{ij}(p^2,\tau)$$
(28)

# High-energy expansion

$$\operatorname{Re}B_{ij}(s,\tau) = -\frac{\tau}{16\pi} \left( 1 - \frac{M_i^2 + M_j^2}{s} \right) + \mathcal{O}(1/s^2) \equiv \operatorname{Re}B_{ij}^{(1)}(s,\tau) + \mathcal{O}(1/s^2)$$
(29)

• For k particles with  $\tau = 1$  and n + 1 - k with  $\tau \neq 1$ 

$$\operatorname{Re}\Sigma(s,\tau) = (1-\tau)\sum_{i,j=0}^{k} a_i a_j B_{ij}^{(1)}(s,1) + \mathcal{O}(1/s^2), \qquad \operatorname{Im}\Sigma(s,\tau) = \mathcal{O}(1/s^2)$$
(30)

# High-energy expansion

$$\operatorname{Re}B_{ij}(s,\tau) = -\frac{\tau}{16\pi} \left( 1 - \frac{M_i^2 + M_j^2}{s} \right) + \mathcal{O}(1/s^2) \equiv \operatorname{Re}B_{ij}^{(1)}(s,\tau) + \mathcal{O}(1/s^2)$$
(29)

• For k particles with  $\tau = 1$  and n + 1 - k with  $\tau \neq 1$ 

$$\operatorname{Re}\Sigma(s,\tau) = (1-\tau)\sum_{i,j=0}^{k} a_i a_j B_{ij}^{(1)}(s,1) + \mathcal{O}(1/s^2), \qquad \operatorname{Im}\Sigma(s,\tau) = \mathcal{O}(1/s^2)$$
(30)

• For  $\tau = 1$  (standard particles and ghosts)

$$\Sigma(s,1) = \mathcal{O}(1/s^2) \tag{31}$$

• For  $\tau = 0$  (k standard particles and n + 1 - k PVP)

$$\operatorname{Re}\Sigma(s,0) = \mathcal{O}(1/s^0), \qquad \operatorname{Im}\Sigma(s,0) = \mathcal{O}(1/s^2)$$
(32)

This difference is crucial in the resummed  $\Omega$  propagator

$$iD(s,\tau) = \frac{1}{N} \frac{iG_r(s,\tilde{\lambda}_i)}{1 - iG_r(s,\tilde{\lambda}_i)\Sigma(s,\tau)}$$
(33)

Explicit example

$$F_1(z) = 1 - z, \qquad G_1(z) = \lambda_0 - \lambda_1 z$$
 (34)

$$\mathcal{L} = \frac{1}{2} (\partial_{\mu} \varphi_a) \left( 1 + \frac{\Box}{M^2} \right) (\partial^{\mu} \varphi_a) - \frac{m^2}{2} \varphi_a \left( 1 + \frac{\Box}{M^2} \right) \varphi_a - \frac{1}{8} \varphi^2 \left( \lambda_0 + \frac{\lambda_1 \Box}{M^2} \right) \varphi^2$$
(35)

Propagators

$$iD_{\varphi}^{ab}(p^2) = -\frac{iM^2\delta^{ab}}{(p^2 - m^2 + i\epsilon)(p^2 - M^2 + i\epsilon)}, \quad \text{with } a_0 = -a_1 = \frac{M^2}{M^2 - m^2}$$
(36)  
$$iD_{\Omega}(p^2) = \lambda_0 - \frac{\lambda_1 p^2}{M^2}$$
(37)

Properties

. . . .

- Superrenormalizable
- For  $\tau = 1$  there are N standard particles of mass m and N ghosts of mass M (violates opt th).
- For  $\tau = 0$  there are N standard particles of mass m and N PVP of mass M (unitary).
- Setting  $F_1(z) = 1$  give a nonrenormalizable theory.



First 40 lines of 3-leg vertex in Stelle gravity. Total is 790 lines, txt file 120 kB



# First 40 lines of 3-leg vertex in Stelle gravity. Total is 790 lines, txt file 120 kB 4-leg vertex txt file is 4 MB

d (mul pul)%d (mu2 pu2)%d (pu3 pu3)%i #Lambda + 2/3%d (mul pul)%d (pu2 pu2)%d (mu3 pu3)%po1 po1%po2 po2%i #Rb + 4/3%d (mu1 pu1)%d (mu2 pu2)%d (mu3 pu3)%po1 po1% pp2,op2\*1 \*Aa + 2/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu2,nu2)\*d (mu3,nu3)\*pp1,op1\*op2,op3\*1 \*Bb + 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pp1,op1\*op2,op3\*1 \*Aa + 2/3\* d (mu1.nu1)\*d (mu2.nu2)\*d (mu3.nu3)\*pp1.pp1\*pp3.pp3\*i \*Bb + 4/3\*d (mu1.nu1)\*d (mu2.nu2)\*d (mu3.nu3)\*pp1.pp1\*pp3.pp3\*i \*Aa - d (mu1.nu1)\*d (mu2.nu2)\* d (mu3,nu3)\*pol,pol\*t \*Cc + 2/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pol,po2\*po3,po3\*t \*Bb + 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pol,po2\*po3,po3\*t \*Aa pp2.pp2+1 +8b + 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pp1.pp3\*pp2.pp2\*1 \*Aa - d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pp1.pp3\*1 \*Cc - d (mu1,nu1)\*d (mu2,nu2)\* <u>d (mu3,nu3)\*pp1.pp3^2\*i \*Aa + 2/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (nu3,nu3)\*pp2.pp3\*i \*Bb + 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pp2.pp2\*pp3.pp3\*i \*Aa</u> 0 (wui,nu1)\*4 (nu2,nu2)\*d (wu3,nu3)\*pp2,pp2\*t \*C - d (wu1,nu1)\*d (wu2,nu2)\*d (wu3,nu3)\*pp2,pp3\*t \*C - d (wu1,nu1)\*d (wu2,nu2)\*d (wu3,nu3)\*pp2,pp3\*t \*Aa d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pp3,pp3\*1 \*Cc + 2\*d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp1(nu3)\*pp1(nu3)\*t \*Cc - 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp1(mu 1\*86 + 8/3\*d (mut out)\*d (mu2 ou2)\*po1(mu3)\*po1(mu3)\*po2 po2\*1 \*aa - d (mut ou1)\*d (mu2 ou2)\*po1(mu3)\*po1(nu3)\*po2 po3\*i \*aa - d (mut ou1)\*d (mu2 ou2)\* pp1(mu3)\*pp1(nu3)\*pp3.pp3\*i \*Aa + d (mu1.nu1)\*d (mu2.nu2)\*pp1(mu3)\*pp2(nu3)\*i \*Cc + 2\*d (mu1.nu1)\*d (mu2.nu2)\*pp1(mu3)\*pp2(nu3)\*pp2(nu3)\*pp2(nu3)\*pp3(nu3)\*p d (mu2,nu2)\*pp1(mu3)\*pp2(nu3)\*pp1,np3\*1 \*Aa + d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp2,np3\*1 \*Aa + d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp3(nu3)\*1 \*Cc + d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp3(nu3)\*pp1,pp3\*i \*Aa - 2/3\*d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp3(nu3)\*pp2,pp2\*i \*Bb - 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\* p3(nu3)\*pp2.pp2\*(\_\*Aa + d (ru1,nu1)\*d (ru2,nu2)\*pp1(nu3)\*pp2(ru3)\*t \*Cc + 2\*d (ru1,nu1)\*d (ru2,nu2)\*pp1(nu3)\*pp2(ru3)\*pp \*pp1(nu3)\*pp2(mu3)\*pp1, pp3\*i \*Aa + d (mu1.nu1)\*d (mu2.nu2)\*pp1(nu3)\*pp2, pp3\*i \*Aa + d (mu1.nu1)\*d (mu2.nu2)\*pp1(nu3)\*pp3(mu3)\*i \*Cc + d (mu1.nu1)\* d\_(mu2, mu2)\*pp1(mu3)\*pp1(mu3)\*pp1,pp31\*\_Ma = /2\*d\_(mu1,mu1)\*d\_(mu2,mu2)\*pp1(mu3)\*pp2(mu3)\*pp2(mu3)\*pp2(mu3)\*pp1(mu3)\*pp1(mu3)\*pp1(mu3)\*pp1(mu3)\*pp3(mu3)\*pp3(mu3)\*pp2(mu3)\*pp2(mu3)\*pp2(mu3)\*pp3(mu3)\*pp d (mu1,nu1)\*d\_(nu2,nu2)\*pp2(nu3)\*p3(nu3)\*t \*Cc · 2/3\*d\_(nu1,nu1)\*d\_(mu2,nu2)\*pp2(nu3)\*pp3(nu3)\*pp1.pp1\*t\*8b · 4/3\*d\_(nu1,nu1)\*d\_(mu2,nu2)\*pp2(nu3)\*p3(nu3)\* pp1.pp1\* \*Aa + d (nu1.nu1)\*d (nu2.nu2)\*pp2(nu3)\*pp3(nu3)\*pp2.pp3\* \*Aa + d (nu1.nu1)\*d (nu2.nu2)\*pp3(nu3)\*i \*Cc - 2/3\*d (nu1.nu1)\*d (nu2.nu2)\* pp2(nu3)\*pp3(nu3)\*pp1.pp1\*1 \*Bb - 4/3\*d (nu1.nu1)\*d (nu2.nu2)\*pp2(nu3)\*pp3.pp1\*1 \*Aa + d (nu1.nu1)\*d (nu2.nu2)\*pp2(nu3)\*pp3(nu3)\*pp1.pp1\*1 \*Aa + d (mu1, nu1)\*d (mu2, nu2)\*pp3(nu3)\*pp3(nu3)\*l \*CC - 2/3\*d (mu1, nu1)\*d (mu2, nu2)\*pp3(nu3)\* pn].pn]+i \*Aa - 2/3\*d (wil.pul)\*d (mi2.piz)\*pn3(mi3)\*pn3( d (mu2,nu2)\*pp3(mu3)\*pp3(nu3)\*pp2,pp2\*1 \*8b - 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*pp3(mu3)\*pp3(nu3)\*pp2,pp2\*1 \*Aa + d (mu1,nu1)\*d (mu2,mu3)\*d (nu2,nu3)\*t \*Lambda - 2/ C(m2,m2) pp2(m3) pp2(m3) pp2(m3) pp2(m3) pp2(m3) m3) = (m2,m3) = Bb - 4/3\*d (mu1,nu1)\*d (mu2,nu3)\*d (nu2,nu3)\*d (nu2,nu pp1.pp2\*pp1.pp3\*1\_\*Aa - 1/2\*d\_(mu1,mu1)\*d\_(mu2,mu3)\*d\_(nu2,nu3)\*pp1.pp2\*pp3.pp3\*1\_\*Aa + d\_(mu1,nu1)\*d\_(mu2,mu3)\*d\_(nu2,nu3)\*pp1.pp2\*1\_\*Cc + d\_(mu1,nu1)\*d\_(mu2,mu3)\*d\_(nu2,mu3 d (mu2,mu3)\*d (nu2,nu3)\*op1.pp2^2\*1 \*Aa - 1/2\*d (mu1,nu1)\*d (mu2,mu3)\*d (nu2,nu3)\*pp1,pp3\*pp2,pp3\*pp2,pp3\*pp2,pp3\*pp1,  $(m_1, m_2) = (m_2, m_3) = (m_$ pp2 pp2+i +(c + 3/2\*d (mu1 pu1)\*d (mu2 mu3)\*d (mu2 mu3)\*nd(mu2 mu3)\*nd(mu2 mu3)\*d (mu2 mu3)\*d (mu2 mu3)\*nd(mu2 mu3)\*n \*noi(nu3)\*i \*Cc + d (nu1.nu1)\*d (nu2.nu3)\*noi(nu pp1(nu2)\*pp3(nu3)\*pp3,pp3\*1 \*Aa - d (mu1,nu1)\*d (mu2,nu3)\*pp1(nu2)\*pp2(nu3)\*1 \*Cc - d (mu1,nu1)\*d (mu2,nu3)\*pp1(nu2)\*pp2(nu3)\*pp1(nu2)\*pp1(nu2)\*pp2(nu3)\*pp1(nu2)\*pp1(nu2)\*pp1(nu2)\*pp1(nu2)\*pp2(nu3)\*pp1(nu2)\*pp1 d (mu2,mu3)\*pp1(nu2)\*pp2(nu3)\*pp3,pp3\*i \*Aa · d (mu1,nu1)\*d (mu2,mu3)\*pp1(nu2)\*pp3(nu3)\*i \*Cc · d (mu1,nu1)\*d (mu2,mu3)\*pp1(nu2)\*pp3(nu3 d (mu1,nu1)\*d (nu2,ru3)\*pp1(nu2)\*pp3(nu3)\*pp1,pp3\*t \*Aa - 1/2\*d (mu1,nu1)\*d (ru2,ru3)\*pp1(nu2)\*pp3(nu3)\*pp2,pp3\*t \*Aa - d (ru1,nu1)\*d (ru2,ru3)\*pp1(nu3)\* pp2/nu2)\*i \*Cc + d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp1 pp(na; ) = C\_C ((na; na)) = (na; na)) = (pp(na; ) = pp(na; ) = (pp(na; na)) = ((na; na)) = (pp(na; ) = (pp(na; na)) = (pp(na; ) = (pp(na; (mu2,mu3)\*pp2(nu2)\*pp2(nu3)\*pp3.pp3\*i \*Aa - d (mu1,nu1)\*d (mu2,mu3)\*pp2(nu2)\*pp3(nu3)\*i\_\*Cc + 2/3\*d (mu1,nu1)\*d (mu2,mu3)\*pp2(nu2)\*pp3(nu3)\*pp1.pp1i\_\*Bb +

First 40 lines of 3-leg vertex in Stelle gravity. Total is 790 lines, txt file 120 kB 4-leg vertex txt file is 4 MB 5-leg vertex txt file is 165 MB

d (mu1 pu1)\*d (mu2 pu2)\*d (mu3 pu3)\*i \*Lambda + 2/3\*d (mu1 pu1)\*d (mu2 pu2)\*d (mu3 pu3)\*po1 po1\*po2 po2\*i \*Bb + 4/3\*d (mu1 pu1)\*d (mu2 pu2)\*d (mu3 pu3)\*po1 po1\* pp2,op2\*1 \*Aa + 2/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu2,nu2)\*d (mu3,nu3)\*pp1,op1\*op2,op3\*1 \*Bb + 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pp1,op1\*op2,op3\*1 \*Aa + 2/3\* d (mu1.nu1)\*d (mu2.nu2)\*d (mu3.nu3)\*pp1.pp1\*pp3.pp3\*i \*Bb + 4/3\*d (mu1.nu1)\*d (mu2.nu2)\*d (mu3.nu3)\*pp1.pp1\*pp3.pp3\*i \*Aa - d (mu1.nu1)\*d (mu2.nu2)\*  $(m_2, m_2) = (m_2, m_2) = (m_$ p2.op2\*1 \*Bb + 4/3\*d (mu1.nu1)\*d (mu2.nu2)\*d (mu3.nu3)\*po1.op3\*p2.op2\*1 \*Aa - d (mu1.nu1)\*d (mu2.nu2)\*d (mu3.nu3)\*po1.op3\*1 \*Cc - d (mu1.nu1)\*d (mu2.nu2)\* <u>d (mu3,nu3)\*pp1.pp3^2\*i \*Aa + 2/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (nu3,nu3)\*pp2.pp3\*i \*Bb + 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pp2.pp2\*pp3.pp3\*i \*Aa</u> 0 (wui,nu1)\*4 (nu2,nu2)\*d (wu3,nu3)\*pp2,pp2\*t \*C - d (wu1,nu1)\*d (wu2,nu2)\*d (wu3,nu3)\*pp2,pp3\*t \*C - d (wu1,nu1)\*d (wu2,nu2)\*d (wu3,nu3)\*pp2,pp3\*t \*Aa d (mu1,nu1)\*d (mu2,nu2)\*d (mu3,nu3)\*pp3,pp3\*i \*CC + 2\*d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp1(nu3)\*i \*CC - 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp1(mu 1\*8b - 8/3\*( (mu1.ou1)\*d (mu2.ou2)\*oo1(mu3)\*oo1(ou3)\*oo2.oo2\*i \*Aa - d (mu1.ou1)\*d (mu2.ou2)\*oo1(mu3)\*oo1(ou3)\*oo2.oo3\*i \*Aa - d (mu1.ou1)\*d (mu2.ou2)\* pp1(mu3)\*pp1(nu3)\*pp3.pp3\*i \*Aa + d (mu1.nu1)\*d (mu2.nu2)\*pp1(mu3)\*pp2(nu3)\*i \*Cc + 2\*d (mu1.nu1)\*d (mu2.nu2)\*pp1(mu3)\*pp2(nu3)\*pp2(nu3)\*pp2(nu3)\*pp3(nu3)\*p d (mu2,nu2)\*pp1(mu3)\*pp2(nu3)\*pp1,np3\*1 \*Aa + d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp2,np3\*1 \*Aa + d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp3(nu3)\*1 \*Cc + d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp3(nu3)\*pp1,pp3\*i \*Aa - 2/3\*d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\*pp3(nu3)\*pp2,pp2\*i \*Bb - 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*pp1(mu3)\* p3(nu3)\*pp2.pp2\*(\_\*Aa + d (ru1,nu1)\*d (ru2,nu2)\*pp1(nu3)\*pp2(ru3)\*t \*Cc + 2\*d (ru1,nu1)\*d (ru2,nu2)\*pp1(nu3)\*pp2(ru3)\*pp2(ru3)\*pp2(ru3)\*pp2(ru3)\*pp1(ru3)\*pp2(ru3)\*pp \*pp1(nu3)\*pp2(mu3)\*pp1, pp3\*i \*Aa + d (mu1.nu1)\*d (mu2.nu2)\*pp1(nu3)\*pp2, pp3\*i \*Aa + d (mu1.nu1)\*d (mu2.nu2)\*pp1(nu3)\*pp3(mu3)\*i \*Cc + d (mu1.nu1)\* d\_(mu2,mu2)\*pp1(mu3)\*pp1(mu3)\*pp1(mu3)\*pp2(mu3)\*pp2(mu3)\*p1(mu3,mu2)\*pp2(mu3)\*pp2(mu pp2(nu3)\*pp2 d (mu1,nu1)\*d\_(nu2,nu2)\*pp2(nu3)\*p3(nu3)\*t \*Cc · 2/3\*d\_(nu1,nu1)\*d\_(mu2,nu2)\*pp2(nu3)\*pp3(nu3)\*pp1.pp1\*t\*8b · 4/3\*d\_(nu1,nu1)\*d\_(mu2,nu2)\*pp2(nu3)\*p3(nu3)\*  $0_{1,001}^{-1,002}$ ,  $100_{2,1002}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,1002}^{-1,002}$ ,  $100_{2,1002}^{-1,002}$ ,  $100_{2,1002}^{-1,002}$ ,  $100_{2,1002}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{2,102}^{-1,002}^{-1,002}$ ,  $100_{2,102}^{-1,002}$ ,  $100_{$ pp2(nu3)\*pp3(nu3)\*pp1.pp1\*(\*Bb - 4/3\*d (nu1.nu1)\*d (nu2.nu2)\*pp2(nu3)\*pp3(nu3)\*pp1.pp1\*(\*Aa + d (nu1.nu1)\*d (nu2.nu2)\*pp2(nu3)\*pp3(nu3)\*pp d (mu1, nu1)\*d (mu2, nu2)\*pp3(nu3)\*pp3(nu3)\*l \*CC - 2/3\*d (mu1, nu1)\*d (mu2, nu2)\*pp3(nu3)\*  $p_{1}$   $p_{1}$   $p_{2}$   $p_{3}$   $p_{3$ d (mu2,nu2)\*pp3(mu3)\*pp3(nu3)\*pp2,pp2\*1 \*8b - 4/3\*d (mu1,nu1)\*d (mu2,nu2)\*pp3(mu3)\*pp3(nu3)\*pp2,pp2\*1 \*Aa + d (mu1,nu1)\*d (mu2,mu3)\*d (nu2,nu3)\*t \*Lambda - 2/ C(m2,m2) pp2(m3) pp2(m3) pp2(m3) pp2(m3) pp2(m3) m3) = (m2,m3) = Bb - 4/3\*d (mu1,nu1)\*d (mu2,nu3)\*d (nu2,nu3)\*d (nu2,nu pp1.pp2\*pp1.pp3\*1\_\*Aa - 1/2\*d\_(mu1,mu1)\*d\_(mu2,mu3)\*d\_(nu2,nu3)\*pp1.pp2\*pp3.pp3\*1\_\*Aa + d\_(mu1,nu1)\*d\_(mu2,mu3)\*d\_(nu2,nu3)\*pp1.pp2\*1\_\*Cc + d\_(mu1,nu1)\*d\_(mu2,mu3)\*d\_(nu2,mu3 d (mu2,mu3)\*d (nu2,nu3)\*pp1.pp2\*2\*i \*Aa - 1/2\*d (mu1,mu1)\*d (mu2,mu3)\*d (nu2,nu3)\*pp1.pp3\*pp2.pp2\*i \*Aa + d (mu1,mu1)\*d (mu2,mu3)\*d (nu2,nu3)\*pp1.pp3\*i \*Cc +  $(m_1, m_2) = (m_2, m_3) = (m_$ pp2 pp2+i +(c + 3/2\*d (mu1 pu1)\*d (mu2 mu3)\*d (mu2 mu3)\*nd(mu2 mu3)\*nd(mu2 mu3)\*d (mu2 mu3)\*d (mu2 mu3)\*nd(mu2 mu3)\*n \*pp1(nu3)\*i\_\*Cc + d (mu1,nu1)\*d (mu2,nu3)\*pp1(nu2)\*pp1(nu3)\*pp2.pp2\*i\_\*Aa + d (mu1,nu1)\*d (mu2,nu3)\*pp1(nu2)\*pp1(nu3)\*pp2.pp3\*i\_\*Aa + d (mu1,nu1)\*d (mu2,mu3)\* pp1(nu2)\*pp3(nu3)\*pp3,pp3\*1 \*Aa - d (mu1,nu1)\*d (mu2,nu3)\*pp1(nu2)\*pp2(nu3)\*1 \*Cc - d (mu1,nu1)\*d (mu2,nu3)\*pp1(nu2)\*pp2(nu3)\*pp1(nu2)\*pp1(nu2)\*pp2(nu3)\*pp1(nu2)\*pp1(nu2)\*pp1(nu2)\*pp1(nu2)\*pp2(nu3)\*pp1(nu2)\*pp1 d (mu2,mu3)\*pp1(nu2)\*pp2(nu3)\*pp3,pp3\*i \*Aa · d (mu1,nu1)\*d (mu2,mu3)\*pp1(nu2)\*pp3(nu3)\*i \*Cc · d (mu1,nu1)\*d (mu2,mu3)\*pp1(nu2)\*pp3(nu3 d (mu1,nu1)\*d (nu2,ru3)\*pp1(nu2)\*pp3(nu3)\*pp1,pp3\*t \*Aa - 1/2\*d (mu1,nu1)\*d (ru2,ru3)\*pp1(nu2)\*pp3(nu3)\*pp2,pp3\*t \*Aa - d (ru1,nu1)\*d (ru2,ru3)\*pp1(nu3)\* pp2/nu2)\*i \*Cc + d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp2/nu2)\*pp1 = 0,0\*i \*Aa + 1/2\*d (mu1, nu1)\*d (mu2, mu3)\*pp1(nu3)\*pp1 pp(na; ) = C\_C ((na; na)) = (na; na)) = (pp(na; ) = pp(na; ) = (pp(na; na)) = ((na; na)) = (pp(na; ) = (pp(na; na)) = (pp(na; ) = (pp(na; (mu2,mu3)\*pp2(nu2)\*pp2(nu3)\*pp3.pp3\*i \*Aa - d (mu1,nu1)\*d (mu2,mu3)\*pp2(nu2)\*pp3(nu3)\*i\_\*Cc + 2/3\*d (mu1,nu1)\*d (mu2,mu3)\*pp2(nu2)\*pp3(nu3)\*pp1.pp1i\_\*Bb +

First 40 lines of 3-leg vertex in Stelle gravity. Total is 790 lines, txt file 120 kB 4-leg vertex txt file is 4 MB 5-leg vertex txt file is 165 MB 6-leg vertex txt file is 5.7 GB.

# Results

# Cross sections

For  $\varphi^a \varphi^a \to \varphi^b \varphi^b$  with  $a \neq b$ 

 $\bullet$  Tree-level cross section

$$\sigma(s) = \frac{a_0^2}{16\pi s} \left(\lambda_0 - \lambda_1 \frac{s}{M^2}\right)^2 \sim \frac{s}{M^4} \text{ at high energies},\tag{38}$$

• Resummed cross sections

$$\sigma(s,\tau) = \frac{a_0^2}{16\pi s} |D(s,\tau)|^2$$
(39)

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(39)

Ghost vs PVP

$$\sigma_{\rm gh}(s) \equiv \sigma(s,1) \sim \frac{\tilde{\lambda}_1^2 s}{16\pi (M^2 - m^2)^2 N^2}, \qquad \sigma_{\rm PVP}(s) \equiv \sigma(s,0) \sim \frac{16\pi (M^2 - m^2)^2}{M^4 N^2} \frac{1}{s}$$
(40)

Nonrenormalizable case

$$F_0(z) = 1, \qquad G_1(z) = \lambda_0 - \lambda_1 z \tag{41}$$

$$\sigma_{\rm nr}(s) \sim \frac{16\pi}{s\left(1 + \ln^2 \frac{s}{m^2}\right)N^2} \tag{42}$$

Cross sections (plot)



# Positivity bounds



- Nonperturbative resummations can cure appartent violations of unitarity bounds in renormalizable theories with PVP, as well as in unitary nonrenormalizable theories
- Renormalizable theories with ghosts cannot be cured in this way
- The resummations add new poles in the cross sections that could be interpreted as resonances

Difficulties with gravity

- Resum quartic vertices
- Justify the resummation (could be just a prescription)
- Long computations...