

High-Energy Behavior of Scattering Amplitudes in Theories with Purely Virtual Particles

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Outline

- ① Introduction on PVP
- ② Scattering in quantum gravity
- ③ Unitarity Bounds and Perturbativity
- ④ $O(N)$ models and large- N expansion
- ⑤ Results

Introduction on PVP

Particle Physics and Quantum Gravity

In particle physics

Locality, renormalizability and unitarity (plus symmetries)



Standard Model

Particle Physics and Quantum Gravity

In particle physics

Locality, renormalizability and unitarity (plus symmetries)



Standard Model

In gravity

Einstein gravity

$$S = \int \sqrt{-g} R$$

Unitary ✓ Renormalizable ✗

Stelle gravity

$$S = \int \sqrt{-g} [R + R^2 + R_{\mu\nu}^2]$$

Unitary ✗ Renormalizable ✓

Degrees of freedom in Stelle theory

$$S_{\text{QG}}(g) = -\frac{1}{16\pi} \int \sqrt{-g} \left(M_{\text{pl}}^2 R - \frac{\xi}{6} R^2 + \frac{\alpha}{2} C_{\mu\nu\rho\sigma} C^{\mu\nu\rho\sigma} \right), \quad g_{\mu\nu} = \eta_{\mu\nu} + \sqrt{32\pi} h_{\mu\nu}.$$

DOF

Massless graviton

Massive scalar ϕ

Massive spin-2 ghost χ

$h_{\mu\nu}$ propagator

$$\frac{i}{p^2} \left[\frac{\Pi_{\mu\nu\rho\sigma}^{(2)}}{(M_{\text{pl}}^2 - \alpha p^2)} - \frac{\Pi_{\mu\nu\rho\sigma}^{(0)}}{2(M_{\text{pl}}^2 - \xi p^2)} \right] + \text{g.f.} = \frac{i}{2M_{\text{pl}}^2} \left[\frac{2\Pi_{\mu\nu\rho\sigma}^{(2)} - \Pi_{\mu\nu\rho\sigma}^{(0)}}{p^2} - \frac{2\Pi_{\mu\nu\rho\sigma}^{(2)}}{p^2 - m_\chi^2} + \frac{\Pi_{\mu\nu\rho\sigma}^{(0)}}{p^2 - m_\phi^2} \right] + \text{g.f.},$$

$$m_\chi^2 = M_{\text{pl}}^2/\alpha, \quad m_\phi^2 = M_{\text{pl}}^2/\xi.$$

Off-shell χ and ϕ are responsible for renormalizability ✓

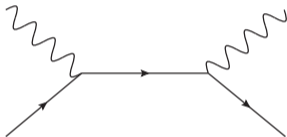
On-shell χ is responsible for the violation of unitarity ✗

Purely Virtual Particles

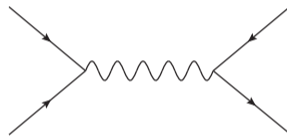
D. Anselmi and MP, JHEP 06 (2017) 066, D. Anselmi JHEP 06 (2017) 086

In quantum field theory a particle can be either "virtual" or "real"(on shell)

Real photon: light



Virtual photon: electromagnetic force

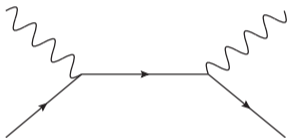


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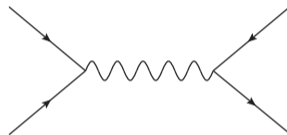
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On-shell and virtual parts are related because of the optical theorem and the Feynman prescription

$$\int d\Pi_f \left| \text{Diagram} \right|^2 = 2 \text{Im} \left[(-i) \text{Diagram} \right]$$

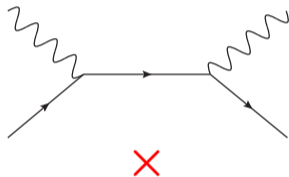
The equation shows the relationship between the squared magnitude of a real particle production amplitude and the imaginary part of a corresponding virtual particle loop diagram. The real particle production amplitude on the left is represented by a diagram with one incoming line and two outgoing lines, where one of the outgoing lines is highlighted in blue. The virtual particle loop diagram on the right is represented by a diagram with two incoming lines and two outgoing lines, where the loop is highlighted in blue.

Purely Virtual Particles

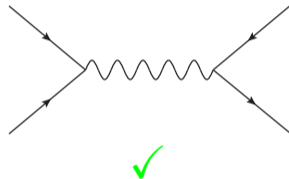
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Purely virtual particles (PVP)

- Only mediate interactions
- Contribute to renormalization
 - Never be on shell
- Consistent with unitarity.

Unitarity

- Optical theorem

$$S = 1 + iT, \quad SS^\dagger = 1 \quad \Leftrightarrow \quad -i(T - T^\dagger) = TT^\dagger.$$

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$$-i(T - T^\dagger) = THT^\dagger, \quad H = \text{diag}(1, \dots, 1, -1, \dots, -1)$$

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- Goal: consistent projection

$$\int d\Pi_f \left| \text{projected away} \right|^2 = 2 \text{Im} \left[(-i) \text{diagram} \right] = \text{diagram} = 0 \text{ by fakeon prescription}$$

Cuts and imaginary parts

PVP cut propagator vanishes. A bubble diagram has only one threshold

$$\begin{array}{c} | \\ | \\ \text{---} \\ | \\ | \end{array} = 0 \Rightarrow \text{---} \begin{array}{c} | \\ | \\ \bigcirc \\ | \\ | \end{array} \text{---} = 0, \quad \text{---} \begin{array}{c} | \\ | \\ \bigcirc \\ | \\ | \end{array} \text{---} = 0.$$

More complicated diagrams have multiple thresholds \Rightarrow Modified functions

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More complicated diagrams have multiple thresholds \Rightarrow Modified functions

$$\text{---} \triangle \overset{\text{---}}{\perp} \neq 0, \quad \text{---} \triangle \overset{\text{---}}{\perp} = 0$$

$$\text{Diagram} \rightarrow \text{Diagram} - \Delta_{\text{Diagram}}^n, \quad n = \# \text{ of PVP inside the loop} \quad (1)$$

For one-loop bubble diagram (this talk): $\Delta_{\text{Bubble}}^{1,2} = \text{Re}(\text{Bubble})$

For triangles and boxes see [A. Melis and MP, PRD 108 \(2023\) 9, 096021](#).

For general procedure [D. Anselmi JHEP 11 \(2021\) 030](#).

Threshold decomposition and spectral identities

D. Anselmi, JHEP 11 (2021) 030

$$G_N = \int \frac{d^D k}{(2\pi)^D} \prod_{a=1}^N \frac{1}{(k - p_a)^2 - m_a^2 + i\epsilon_a} = \int \frac{d^{D-1} \mathbf{k}}{(2\pi)^{D-1}} \left(\prod_{a=1}^N \frac{1}{2\omega_a} \right) G_N^s$$

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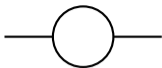
Skeleton diagram

$$G_N^s = \int \frac{dk^0}{2\pi} \prod_{a=1}^N \frac{2\omega_a}{(k^0 - e_a)^2 - \omega_a^2 + i\epsilon_a}, \quad e_a = p_a^0, \quad \omega_a = \sqrt{(\mathbf{k} - \mathbf{p}_a)^2 + m_a^2}$$

Spectral identities

$$G^s + (G^s)^* + \sum_{\text{cuts}} G_c^s = 0.$$

which holds threshold by thresholds.

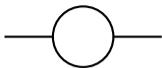


$$B^s = -\frac{i}{e_1 - e_2 - \omega_1 - \omega_2 + i\epsilon} - \frac{i}{e_2 - e_1 - \omega_1 - \omega_2 + i\epsilon}.$$

Using $\frac{i}{x+i\epsilon} = \mathcal{P}\frac{i}{x} + \pi\delta(x)$ everywhere

$$B^s = -i\mathcal{P}_2 - \Delta^{12} - \Delta^{21}$$

$$\mathcal{P}^{ab} = \mathcal{P}\frac{1}{e_a - e_b - \omega_a - \omega_b}, \quad \mathcal{P}_2 = \mathcal{P}^{ab} + \mathcal{P}^{ba}, \quad \Delta^{ab} = \pi\delta(e_a - e_b - \omega_a - \omega_b)$$







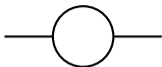
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Diag. \ Terms				
\mathcal{P}_2	$-i$	i	0	0
Δ^{12}	-1	-1	2	0
Δ^{21}	-1	-1	0	2



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If we want to include PVP, we kill all the Δ 's that contain at least one PVP frequency.

Diag.								
\mathcal{P}_3	$-i$	i	0	0	0	0	0	0
$\Delta^{12} Q^{13}$	-1	-1	2	0	0	0	0	0
$\Delta^{23} Q^{21}$	-1	-1	0	2	0	0	0	0
$\Delta^{31} Q^{32}$	-1	-1	0	0	2	0	0	0
$\Delta^{21} Q^{23}$	-1	-1	0	0	0	2	0	0
$\Delta^{32} Q^{31}$	-1	-1	0	0	0	0	2	0
$\Delta^{13} Q^{12}$	-1	-1	0	0	0	0	0	2
$\Delta^{12} \Delta^{13}$	i	$-i$	$2i$	0	0	0	0	$-2i$
$\Delta^{23} \Delta^{21}$	i	$-i$	0	$2i$	0	$-2i$	0	0
$\Delta^{31} \Delta^{32}$	i	$-i$	0	0	$2i$	0	$-2i$	0
$\Delta^{21} \Delta^{31}$	i	$-i$	0	0	$2i$	$-2i$	0	0
$\Delta^{32} \Delta^{12}$	i	$-i$	$2i$	0	0	0	$-2i$	0
$\Delta^{13} \Delta^{23}$	i	$-i$	0	$2i$	0	0	0	$-2i$

$$\mathcal{P}_3 = \mathcal{P}^{12} \mathcal{P}^{13} + \text{cycl} + (e \rightarrow -e), \quad Q^{ab} = \mathcal{P}^{ab} - \mathcal{P} \frac{1}{e_a - e_b - \omega_a + \omega_b}.$$

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Scattering in quantum gravity

Graviton scattering in quantum gravity

In Einstein gravity (unitary, nonrenormalizable)

$$S_{\text{H}}(g) = -\frac{1}{2\kappa^2} \int \sqrt{-g} R, \quad \kappa^2 = 8\pi G, \quad g_{\mu\nu} = \eta_{\mu\nu} + 2\kappa h_{\mu\nu} \quad (2)$$

Tree level amplitudes

$$\mathcal{A}_{hh \rightarrow hh}^{\text{EG}} \sim \kappa^2 s, \quad s = \text{c.o.m. energy squared} \quad (3)$$

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Tree level amplitudes

P. Donà, S. Giaccari, L. Modesto, L. Rachwal, Y. Zhu JHEP 08 (2015) 038.

$$\mathcal{A}_{hh \rightarrow hh}^{\text{SG}} = \mathcal{A}_{hh \rightarrow hh}^{\text{EG}} \quad (5)$$

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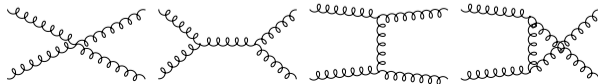
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Equivalent formulation of Stelle gravity

D. Anselmi and MP, JHEP 11 (2018) 021.

$$S_{\text{SG}}(g) = -\frac{1}{2\kappa^2} \int \sqrt{-g} \left[\zeta R - \frac{1}{6\xi} R^2 + \frac{1}{2\alpha} C_{\mu\nu\rho\sigma} C^{\mu\nu\rho\sigma} \right] \quad (6)$$

Equivalent action

auxiliary fields ϕ , $\chi_{\mu\nu}$ + Weyl transformation + field redefinitions.

↓

$$S(g, \phi, \chi) = -\frac{1}{2\kappa^2} \int \sqrt{-g} \left[\zeta R + \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right] + S_{\chi\chi}(g, \chi) + S_{\text{int}}(g, \chi, \phi) \quad (7)$$

$$V(\phi) = \frac{m_\phi^2}{\kappa^2} (1 - e^{\tilde{\kappa}\phi})^2, \quad S_{\chi\chi} = -S_{\text{Pauli-Fierz}}$$

Vertices in S_{int}



$$h\chi\chi, \quad h\phi\phi, \quad ,hh\chi\chi, \quad hh\phi\phi, \quad \chi\phi\phi, \quad \chi\chi\phi\phi \quad \text{etc...} \quad (8)$$



$$hh\chi, \quad hh\phi, \quad \text{etc..} \quad (9)$$

Unitarity Bounds and Perturbativity

Unitarity bounds and perturbativity

Usual derivation of Unitarity bounds for 2-to-2 scattering of scalars

$$\sigma(s) = \frac{1}{32\pi s} \int_{-1}^1 dv |\mathcal{M}(s, v)|^2, \quad v \equiv \cos \theta. \quad (10)$$

$$S = 1 + iT, \quad S^\dagger S = 1, \quad -i(T - T^\dagger) = \frac{1}{2} T^\dagger T \quad (11)$$

↓

$$\text{Im} \mathcal{M}(s, 1) = \sqrt{\kappa(s, m_1^2, m_2^2)} \sum_X \sigma_X(s) \geq \sqrt{\kappa(s, m_1^2, m_2^2)} \sigma(s), \quad (12)$$

$$\kappa(x, y, z) = x^2 + y^2 + z^2 - xy - xz - yz \quad (13)$$

Partial wave expansion in Legendre polynomials

$$\mathcal{M}(s, v) = 16\pi \sum_{j=0}^{\infty} (2j+1) A_j(s) P_j(v), \quad \sigma(s) = \frac{16\pi}{s} \sum_{j=0}^{\infty} (2j+1) |A_j(s)|^2 \quad (14)$$

Unitarity bounds and perturbativity

$$\sum_{j=0}^{\infty} (2j+1) \operatorname{Im} A_j(s) \geq \frac{\sqrt{\kappa(s, m_1^2, m_2^2)}}{s} \sum_{j=0}^{\infty} (2j+1) |A_j(s)|^2. \quad (15)$$

For elastic scattering only we have the equality, which leads to

$$|\mathcal{A}_j| \leq 1, \quad 0 \leq \operatorname{Im} \mathcal{A}_j \leq 1, \quad |\operatorname{Re} \mathcal{A}_j| \leq \frac{1}{2}, \quad \mathcal{A}_j(s) = \frac{\sqrt{\kappa(s, m_1^2, m_2^2)}}{s} A_j(s) \quad (16)$$

So far Unitarity \Rightarrow (15). Therefore, if (15) is violated \Rightarrow Unitarity is violated.

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• Typical argument:

$$\mathcal{M} = \mathcal{M}^{\text{tree}} + \mathcal{M}^{\text{loops}}, \quad \mathcal{M}^{\text{tree}} = 16\pi A_0^{\text{tree}} + \dots, \quad A_0^{\text{tree}} \sim s\text{-channel diagrams} \quad (17)$$

✗ if $\mathcal{A}_0^{\text{tree}}$ violates (16) at some scale, then the theory violates unitarity at some scale.

Unitarity **and** perturbativity ($|\mathcal{M}^{\text{loop}}| < |\mathcal{M}^{\text{tree}}|$) \Rightarrow (16) for $\mathcal{A}_0^{\text{tree}}$

✓ if $\mathcal{A}_0^{\text{tree}}$ violates (16) at some scale, then **either** unitarity or perturbativity is violated at some scale

Diagrammatic optical theorem

Unitarity can be checked by means of the cutting equations

$$2\text{Im}(-iG) = -\sum_c G_c \quad (18)$$

which hold for any diagram G for any local QFT.

- Absence of ghosts \Rightarrow (18) is the diagrammatic version of the unitarity equation.
- The cutting equations (18) relies only on the diagrammatic expansion.
- No assumptions on the behavior of G .

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- No assumptions on the behavior of G .

\Downarrow

If $\mathcal{A}_0^{\text{tree}}$ violates $|\mathcal{A}_0^{\text{tree}}| \leq 1$ etc... at some scale, then perturbativity is violated at some scale.

$O(N)$ models

$O(N)$ models and large- N diagrammatics

$$\mathcal{L}(\varphi) = \frac{1}{2} \partial_\mu \varphi^a \partial^\mu \varphi^a - \frac{1}{2} m^2 \varphi^a \varphi^a - \frac{g}{8} (\varphi^a \varphi^a)^2, \quad a = 1, \dots, N \quad (19)$$

Identify the orders in N

$$-\frac{g}{8} (\varphi^a \varphi^a)^2 \rightarrow \frac{1}{2g} \Omega^2 - \frac{1}{2} \Omega \varphi^a \varphi^a, \quad \tilde{g} = gN \quad (\text{'t Hooft coupling}) \quad (20)$$

Study the limit $N \rightarrow \infty$ by keeping \tilde{g} fixed.

- Each Ω internal line gives $1/N$
- Each closed φ loop gives N

$O(N)$ models and large- N diagrammatics

$$\mathcal{L}(\varphi) = \frac{1}{2} \partial_\mu \varphi^a \partial^\mu \varphi^a - \frac{1}{2} m^2 \varphi^a \varphi^a - \frac{g}{8} (\varphi^a \varphi^a)^2, \quad a = 1, \dots, N \quad (19)$$

Identify the orders in N

$$-\frac{g}{8} (\varphi^a \varphi^a)^2 \rightarrow \frac{1}{2g} \Omega^2 - \frac{1}{2} \Omega \varphi^a \varphi^a, \quad \tilde{g} = gN \quad (\text{'t Hooft coupling}) \quad (20)$$

Study the limit $N \rightarrow \infty$ by keeping \tilde{g} fixed.

- Each Ω internal line gives $1/N$
- Each closed φ loop gives N

Consider the scattering of $\varphi^a \varphi^a \rightarrow \varphi^b \varphi^b$ with $a \neq b$



All the bubble insertions are order $1/N$ and can be resummed

Higher-derivative $O(N)$ models

$$\mathcal{L}(\varphi) = \frac{1}{2} \partial_\mu \varphi^a F_n \left(\frac{-\square}{M^2} \right) \partial^\mu \varphi^a - \frac{1}{2} m^2 \varphi^a F_n \left(\frac{-\square}{M^2} \right) \varphi^a - \frac{1}{8} \varphi^2 G_r \left(\frac{-\square}{M^2} \right) \varphi^2, \quad \varphi^2 \equiv \varphi^a \varphi^a, \quad (21)$$

$$F_n(z) = 1 + \sum_{i=1}^n f_i z^i, \quad G_r(z) = \sum_{i=0}^r \lambda_i z^i, \quad (22)$$

- Auxiliary fields and 't Hooft couplings

$$-\frac{1}{8} \varphi^2 G_r \left(\frac{-\square}{M^2} \right) \varphi^2 \rightarrow \frac{1}{2} \Omega G_r^{-1} \left(\frac{-\square}{M^2} \right) \Omega - \frac{1}{2} \Omega \varphi^2, \quad \tilde{\lambda}_i = \lambda_i N \quad (23)$$

- Propagators

$$iD_\Omega(p^2) = iG_r(p^2), \quad iD_\varphi^{ab}(p) = \frac{i\delta^{ab}}{(p^2 - m^2 + i\epsilon)F_n(p^2/M^2)} \equiv i\delta^{ab} D_{\text{HD}}(p^2) \quad (24)$$

Bubble diagrams

$$D_{\text{HD}}(p^2) = \frac{a_0}{p^2 - m^2 + i\epsilon} + \sum_{i=1}^n \frac{a_i}{p^2 - M_i^2 + i\epsilon}, \quad \text{with} \quad \sum_{i=0}^n a_i = 0 \quad (25)$$

- Generic bubble diagram

$$B_{ij}(p^2) \equiv \int \frac{d^D q}{(2\pi)^D} \frac{1}{(p+q)^2 - M_i^2 + i\epsilon} \frac{1}{q^2 - M_j^2 + i\epsilon}, \quad (26)$$

$$B_{ij}(p^2, \tau) = \tau \text{Re} B_{ij}(p^2) + i \text{Im} B_{ij}(p^2), \quad (\tau = 0 \text{ for PVP}) \quad (27)$$

- Ω self energy is $N\Sigma(p^2, \tau)$

$$\Sigma(p^2, \tau) = \int \frac{d^D q}{(2\pi)^D} D_{\text{HD}}(p+q) D_{\text{HD}}(q) = \sum_{i,j=0}^n a_i a_j B_{ij}(p^2, \tau) \quad (28)$$

High-energy expansion

$$\operatorname{Re}B_{ij}(s, \tau) = -\frac{\tau}{16\pi} \left(1 - \frac{M_i^2 + M_j^2}{s} \right) + \mathcal{O}(1/s^2) \equiv \operatorname{Re}B_{ij}^{(1)}(s, \tau) + \mathcal{O}(1/s^2) \quad (29)$$

- For k particles with $\tau = 1$ and $n + 1 - k$ with $\tau \neq 1$

$$\operatorname{Re}\Sigma(s, \tau) = (1 - \tau) \sum_{i,j=0}^k a_i a_j B_{ij}^{(1)}(s, 1) + \mathcal{O}(1/s^2), \quad \operatorname{Im}\Sigma(s, \tau) = \mathcal{O}(1/s^2) \quad (30)$$

High-energy expansion

$$\operatorname{Re}B_{ij}(s, \tau) = -\frac{\tau}{16\pi} \left(1 - \frac{M_i^2 + M_j^2}{s} \right) + \mathcal{O}(1/s^2) \equiv \operatorname{Re}B_{ij}^{(1)}(s, \tau) + \mathcal{O}(1/s^2) \quad (29)$$

- For k particles with $\tau = 1$ and $n + 1 - k$ with $\tau \neq 1$

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- For $\tau = 1$ (standard particles and ghosts)

$$\Sigma(s, 1) = \mathcal{O}(1/s^2) \quad (31)$$

- For $\tau = 0$ (k standard particles and $n + 1 - k$ PVP)

$$\operatorname{Re}\Sigma(s, 0) = \mathcal{O}(1/s^0), \quad \operatorname{Im}\Sigma(s, 0) = \mathcal{O}(1/s^2) \quad (32)$$

This difference is crucial in the resummed Ω propagator

$$iD(s, \tau) = \frac{1}{N} \frac{iG_r(s, \tilde{\lambda}_i)}{1 - iG_r(s, \tilde{\lambda}_i)\Sigma(s, \tau)} \quad (33)$$

Explicit example

$$F_1(z) = 1 - z, \quad G_1(z) = \lambda_0 - \lambda_1 z \quad (34)$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \varphi_a) \left(1 + \frac{\square}{M^2}\right) (\partial^\mu \varphi_a) - \frac{m^2}{2} \varphi_a \left(1 + \frac{\square}{M^2}\right) \varphi_a - \frac{1}{8} \varphi^2 \left(\lambda_0 + \frac{\lambda_1 \square}{M^2}\right) \varphi^2 \quad (35)$$

Propagators

$$iD_\varphi^{ab}(p^2) = -\frac{iM^2 \delta^{ab}}{(p^2 - m^2 + i\epsilon)(p^2 - M^2 + i\epsilon)}, \quad \text{with } a_0 = -a_1 = \frac{M^2}{M^2 - m^2} \quad (36)$$

$$iD_\Omega(p^2) = \lambda_0 - \frac{\lambda_1 p^2}{M^2} \quad (37)$$

Properties

- Superrenormalizable
- For $\tau = 1$ there are N standard particles of mass m and N ghosts of mass M (violates opt th).
- For $\tau = 0$ there are N standard particles of mass m and N PVP of mass M (unitary).
- Setting $F_1(z) = 1$ give a nonrenormalizable theory.

Why not doing directly Stelle gravity instead of toy models?

Results

Cross sections

For $\varphi^a \varphi^a \rightarrow \varphi^b \varphi^b$ with $a \neq b$

- Tree-level cross section

$$\sigma(s) = \frac{a_0^2}{16\pi s} \left(\lambda_0 - \lambda_1 \frac{s}{M^2} \right)^2 \sim \frac{s}{M^4} \text{ at high energies,} \quad (38)$$

- Resummed cross sections

$$\sigma(s, \tau) = \frac{a_0^2}{16\pi s} |D(s, \tau)|^2 \quad (39)$$

Cross sections

For $\varphi^a \varphi^a \rightarrow \varphi^b \varphi^b$ with $a \neq b$

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- Resummed cross sections

$$\sigma(s, \tau) = \frac{a_0^2}{16\pi s} |D(s, \tau)|^2 \quad (39)$$

Ghost vs PVP

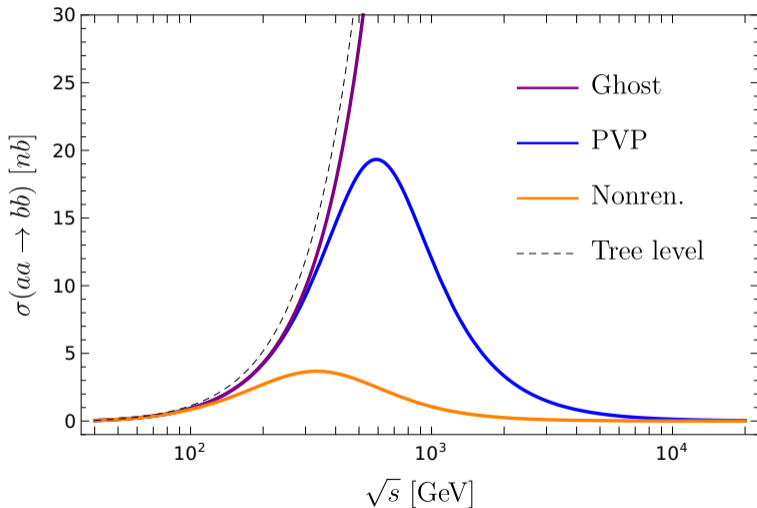
$$\sigma_{\text{gh}}(s) \equiv \sigma(s, 1) \sim \frac{\tilde{\lambda}_1^2 s}{16\pi(M^2 - m^2)^2 N^2}, \quad \sigma_{\text{PVP}}(s) \equiv \sigma(s, 0) \sim \frac{16\pi(M^2 - m^2)^2}{M^4 N^2} \frac{1}{s} \quad (40)$$

Nonrenormalizable case

$$F_0(z) = 1, \quad G_1(z) = \lambda_0 - \lambda_1 z \quad (41)$$

$$\sigma_{\text{nr}}(s) \sim \frac{16\pi}{s \left(1 + \ln^2 \frac{s}{m^2} \right) N^2} \quad (42)$$

Cross sections (plot)

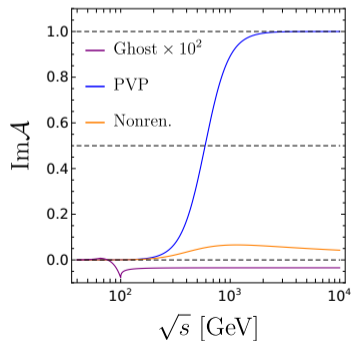
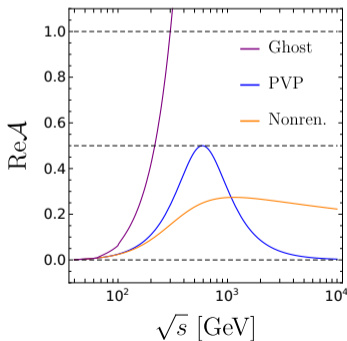
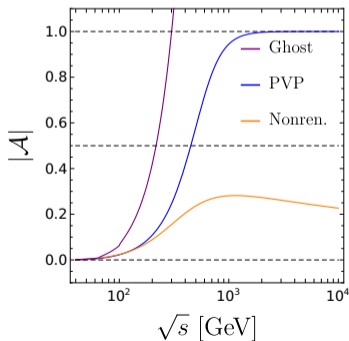


$$m = 15 \text{ GeV}, \quad M = 50 \text{ GeV}, \quad \lambda_0 = 0.3, \quad \lambda_1 = 0.1$$

Positivity bounds

$$\mathcal{A}_{\text{gh}}(s) = -\frac{a_0^2}{16\pi} \left[\sqrt{1 - \frac{4m^2}{s}} + 2\sqrt{\frac{\kappa(s, m^2, M^2)}{s^2}} + \sqrt{1 - \frac{4M^2}{s}} \right] D(s, 1) \quad (43)$$

$$\mathcal{A}_{\text{PVP}}(s) = -\frac{a_0^2}{16\pi} \sqrt{1 - \frac{4m^2}{s}} D(s, 0), \quad \mathcal{A}_{\text{nr}}(s) = -\frac{1}{16\pi} \sqrt{1 - \frac{4m^2}{s}} D_{\text{nr}}(s) \quad (44)$$



Conclusions

- Nonperturbative resummations can cure apparent violations of unitarity bounds in renormalizable theories with PVP, as well as in unitary nonrenormalizable theories
- Renormalizable theories with ghosts cannot be cured in this way
- The resummations add new poles in the cross sections that could be interpreted as resonances

Difficulties with gravity

- Resum quartic vertices
- Justify the resummation (could be just a prescription)
- Long computations...